

Maxwell and His Equations

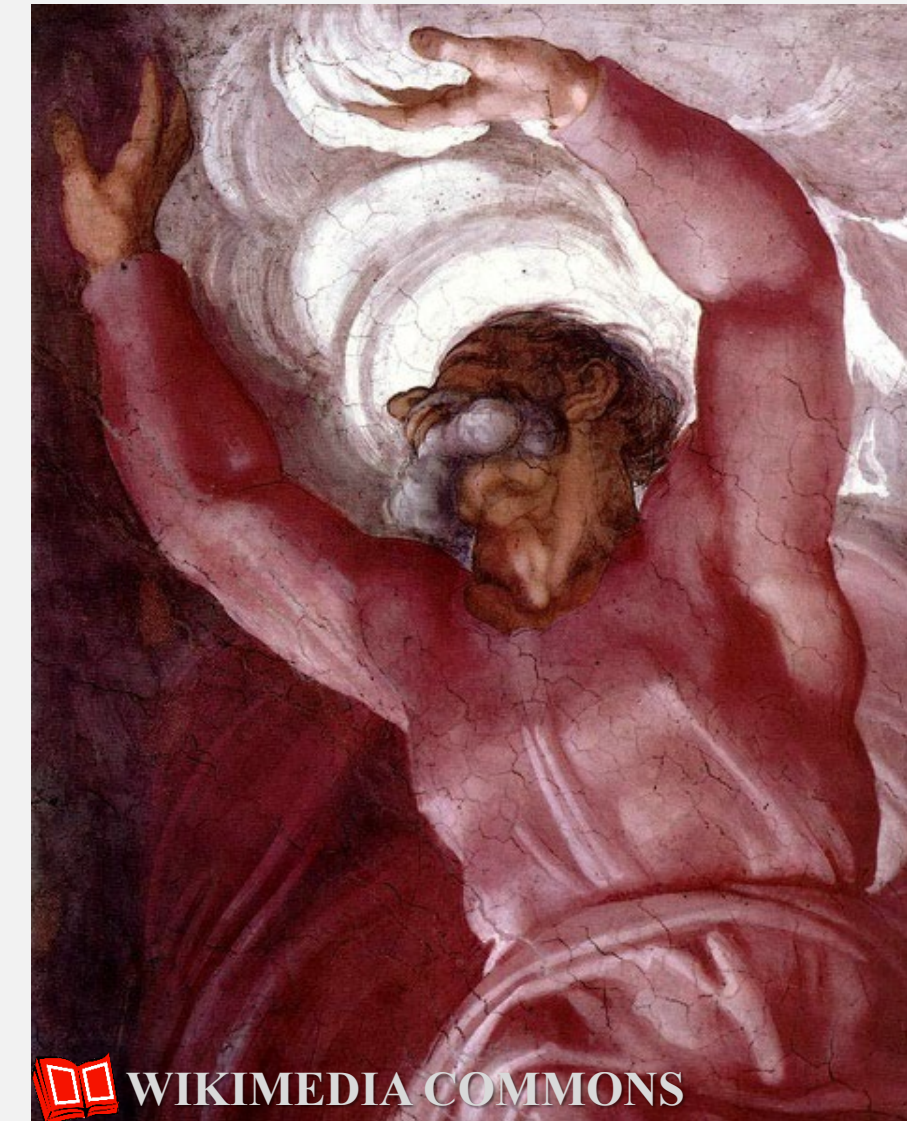
鄭士康 (Shyh-Kang Jeng)

Department of Electrical Engineering

Graduate Institute of Communication Engineering

The Beginning, Genesis 1, Old Testament

1. In the beginning God created the heavens and the earth.
2. Now the earth was formless and empty, darkness was over the surface of the deep, and the Spirit of God was hovering over the waters.
3. And God said, "Let there be light," and there was light.
4. God saw that the light was good, and he separated the light from the darkness.



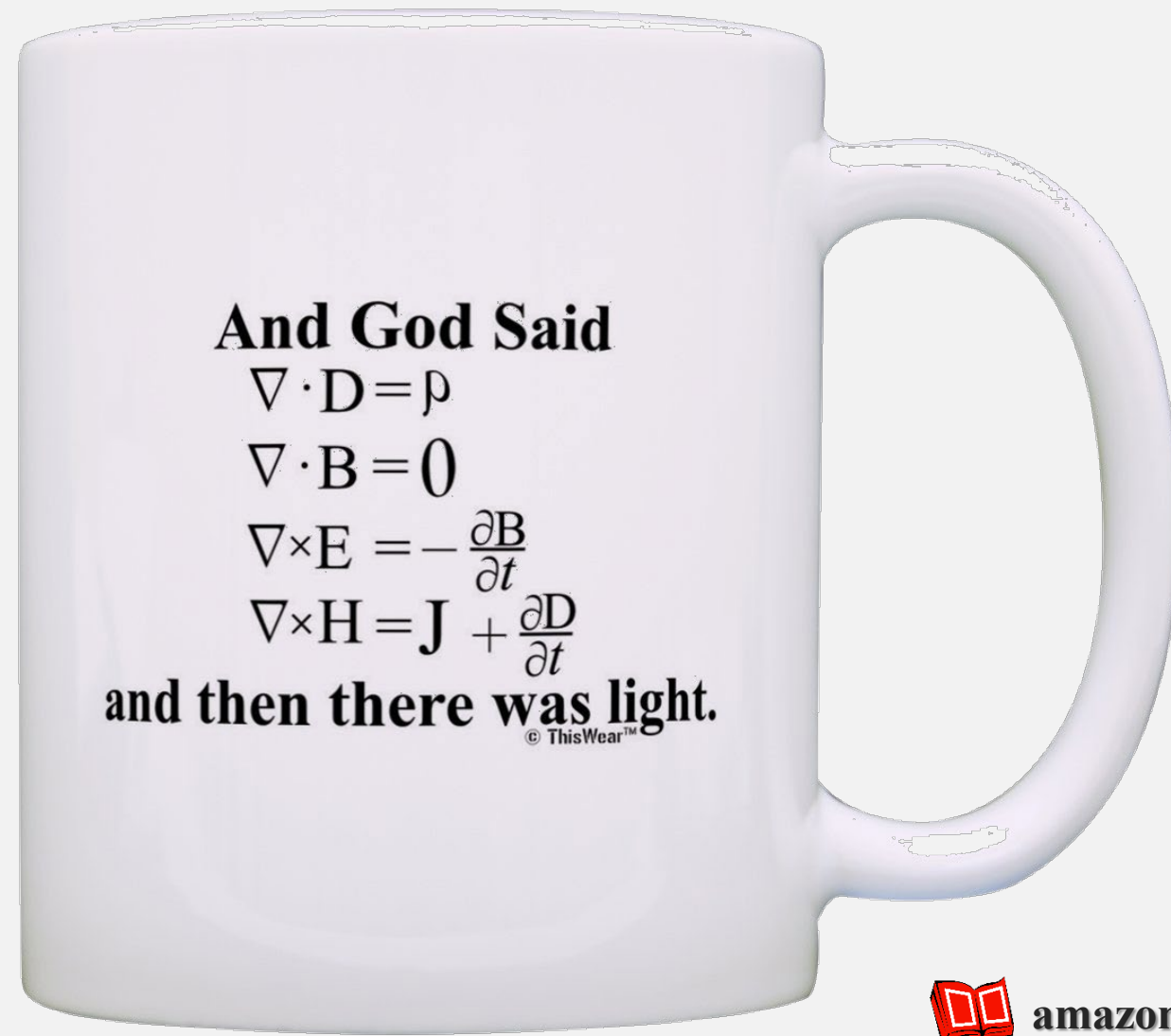
The Sistine Chapel Ceiling / Michelangelo Buonarroti

https://commons.wikimedia.org/wiki/File:Dividing_Light_from_Darkness.jpg



amazon

<https://www.amazon.com/Maxwell-Equations-There-Light-Shirt/dp/B07KVSVK8G>



amazon

https://www.amazon.com/Funny-Maxwells-Equations-There-Coffee/dp/B017OE2KTI/ref=sr_1_1?crd=P8CII24NYLH&keywords=maxwell%2Bequation%2Bmug&qid=1566367884&s=gateway&sprefix=Maxwell%2Beq%2Caps%2C358&sr=8-1&th=1



flickr

<https://flic.kr/p/QGakzx>

Inaugural Lecture of Maxwell, King's College, 1860

In his inaugural lecture, as Professor of Physics and Astronomy at King's College, London, he said:

In this class I hope you will learn not merely results, or formulae applicable to cases that may possibly occur in our practice afterwards, but the principles on which those formulae depend, and without which the formulae are mere mental rubbish.

 **Google Books**

<https://books.google.com.tw/books?id=IMl7AgAAQBAJ&pg=PA8&lpg=PA8&dq=in+this+class+i+hope+you+will+learn+not+merely+results+maxwell&source=bl&ots=UX09wsz5Ux&sig=ACfU3U2-n1ZHGORdIF8DB8ct9LBdYLI-CA&hl=en&sa=X&ved=2ahUKEwjDyOvg-cn1AhUyLqYKHYQgCVcQ6AF6BAgCEAM#v=onepage&q&f=false>



[https://commons.wikimedia.org/wiki/File:James Clerk Maxwell.png](https://commons.wikimedia.org/wiki/File:James_Clerk_Maxwell.png)

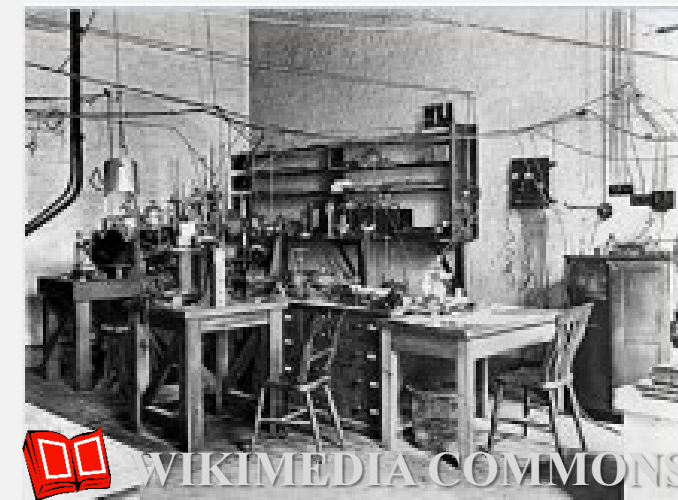
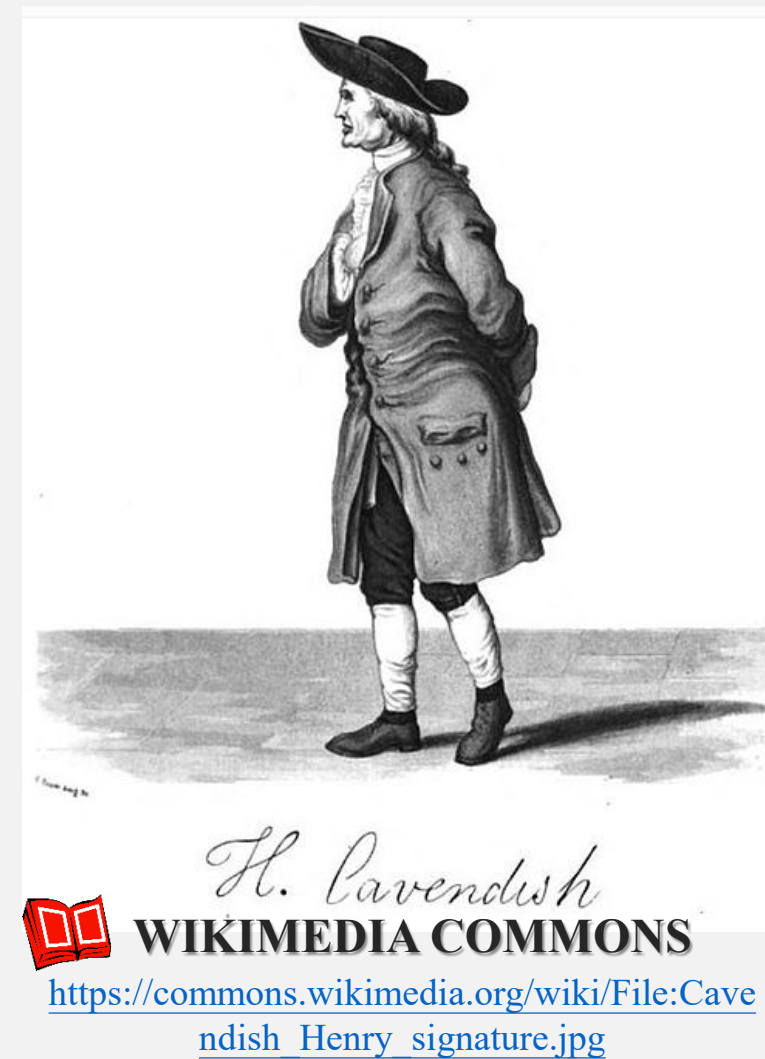
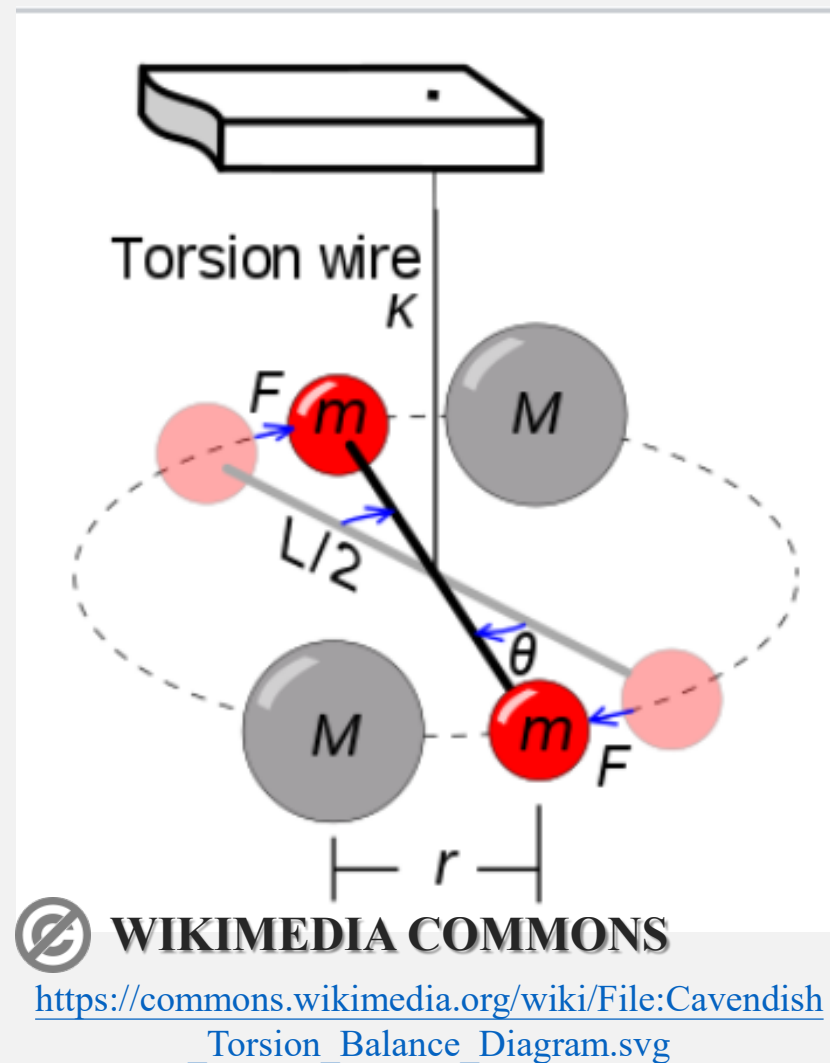
 **WIKIMEDIA COMMONS**



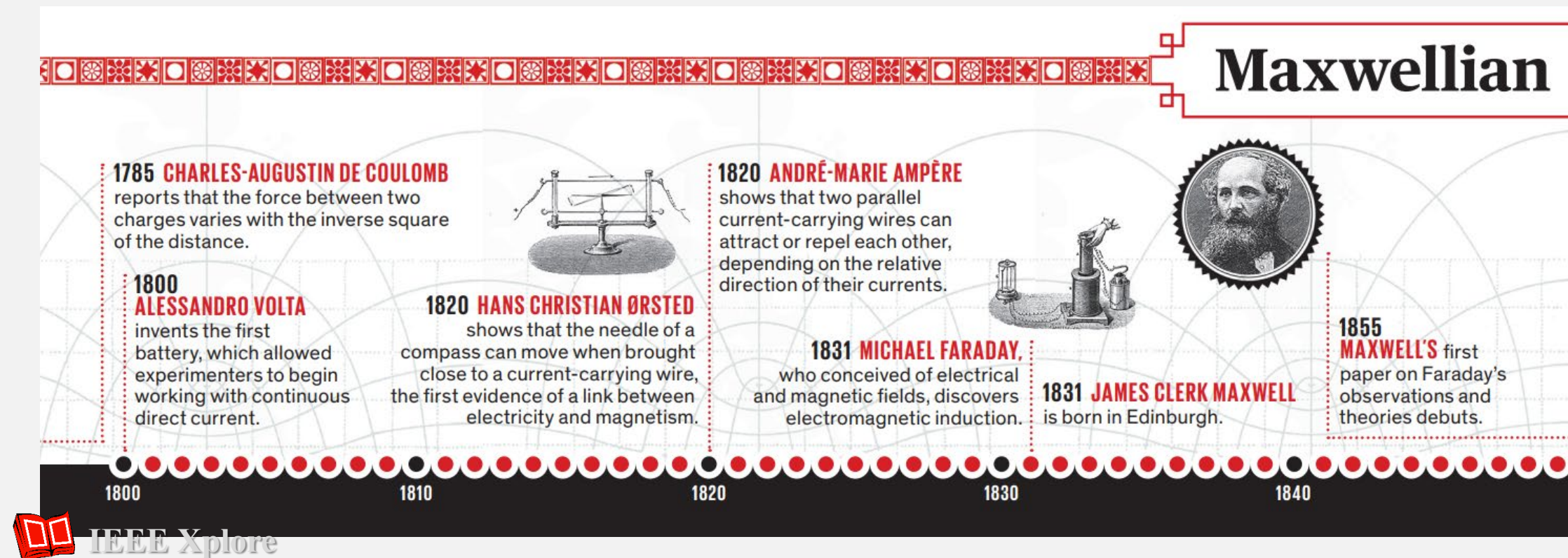
鄭士康 (Shyh-Kang Jeng)

Shyh-Kang Jeng in front of King's College, London, 2018

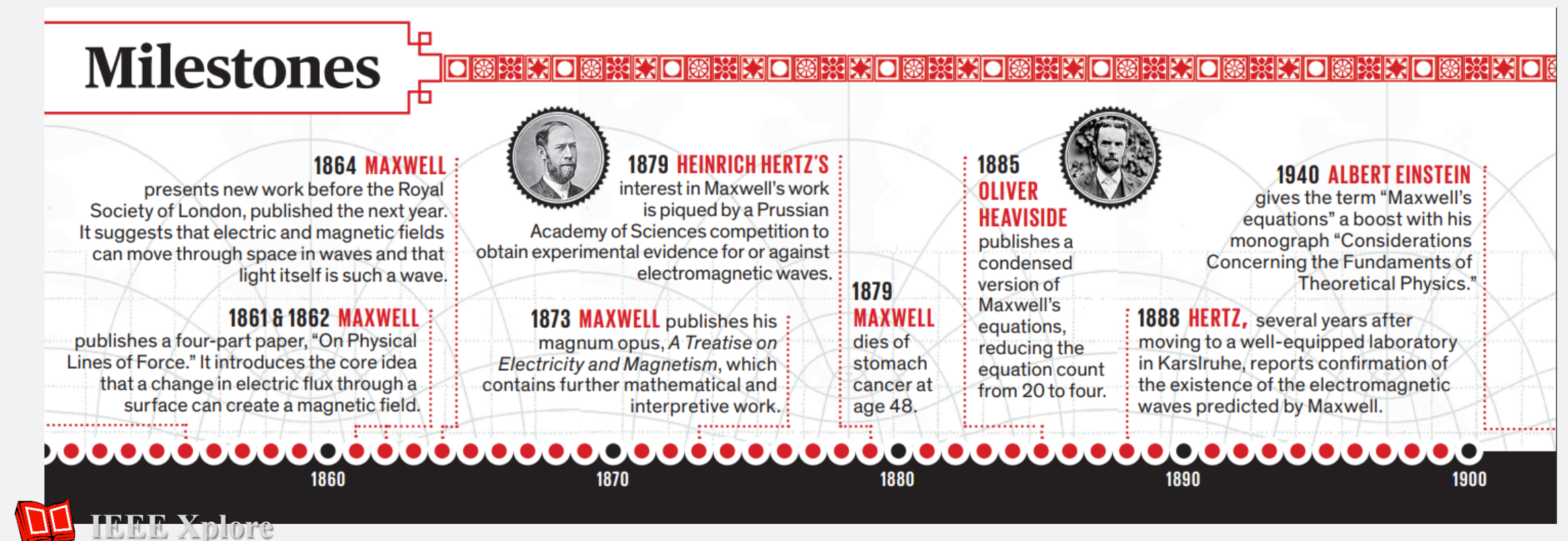
Henry Cavendish and Maxwell Center, Cavendish Laboratory



Long Road to Maxwell's Equations



<https://ieeexplore.ieee.org/stamp/stamp.jsp?tp=&arnumber=6964925>



Outline

- Life of James Clark Maxwell
- Electricity and Magnetism before Maxwell
- Evolution of Maxwell's Equations
- “Derivation” of Maxwell's Equations

Outline

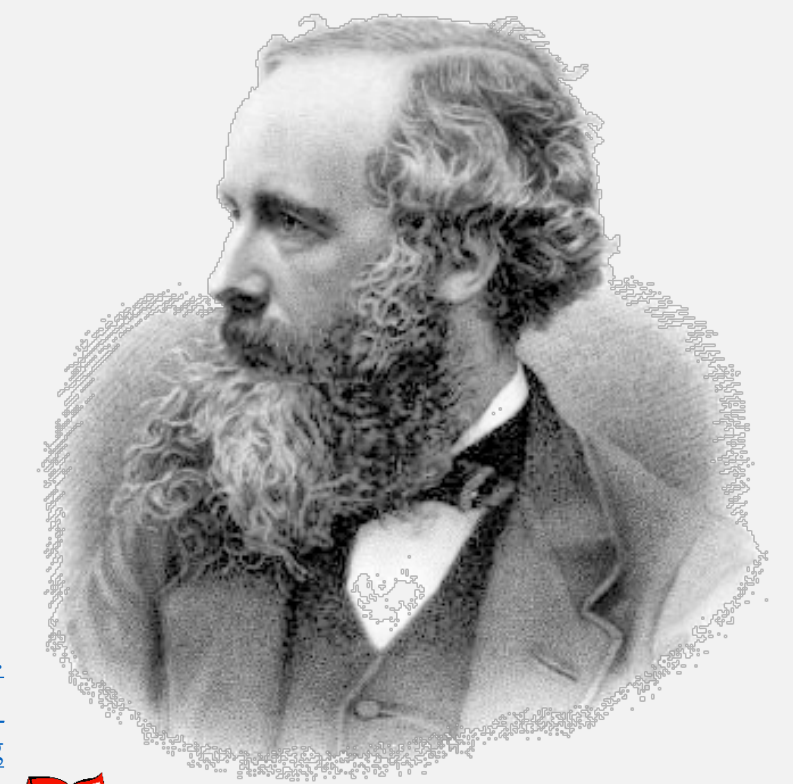
- Life of James Clark Maxwell
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Time Line

DATE	EVENT
1831	Born at 14 India Street, Edinburgh and grew up at Glenlair
1841 - 1847	Edinburgh Academy
1847 - 1850	Edinburgh University
1850 – 1856	Cambridge University
1856 - 1860	Marischal College, Aberdeen
1860 - 1865	King's College, London
1865 - 1871	Glenlair
1871 – 1879	Cambridge University

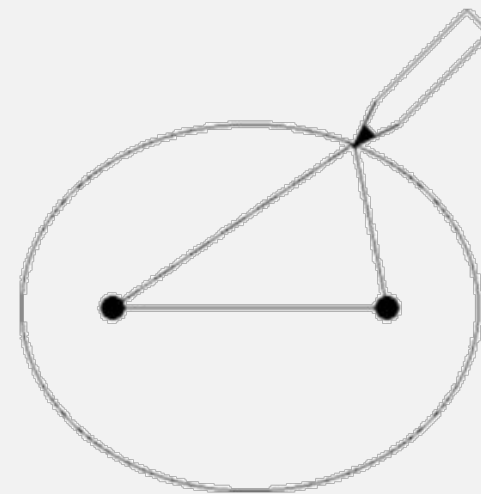
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[https://commons.wikimedia.org/wiki/File:James Clerk Maxwell.png](https://commons.wikimedia.org/wiki/File:James_Clerk_Maxwell.png)



 WIKIMEDIA COMMONS

First Paper, Drawing Oval Curves, 1846



[From the *Proceedings of the Royal Society of Edinburgh*, Vol. II. April, 1846.]

I. On the Description of Oval Curves, and those having a plurality of Foci; with remarks by Professor Forbes. Communicated by PROFESSOR FORBES.

MR CLERK MAXWELL ingeniously suggests the extension of the common theory of the foci of the conic sections to curves of a higher degree of complication in the following manner:—

(1) As in the ellipse and hyperbola, any point in the curve has the sum or difference of two lines drawn from two points or foci = a constant quantity, so the author infers, that curves to a certain degree analogous, may be described and determined by the condition that the simple distance from one focus plus a multiple distance from the other, may be = a constant quantity; or more generally, m times the one distance + n times the other = constant.

(2) The author devised a simple mechanical means, by the wrapping of a thread round pins, for producing these curves. See Figs. 1 and 2. He

Fig. 1. Two Foci. Ratios 1, 2.

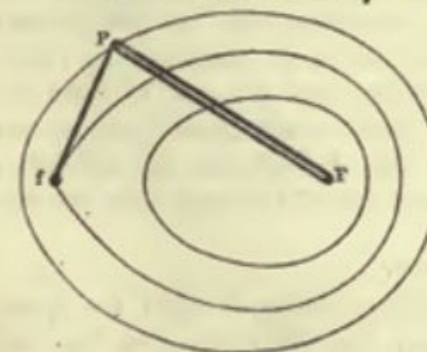
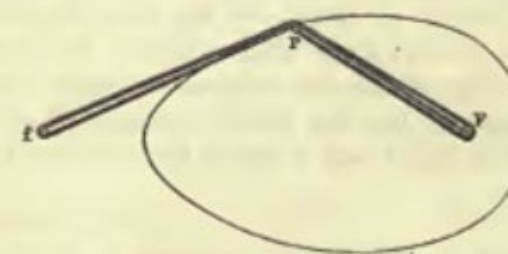


Fig. 2. Two Foci. Ratios 2, 3.



then thought of extending the principle to other curves, whose property should be, that the sum of the simple or multiple distances of any point of

VOL. I.

1



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1854 Smith's Prize Examination

8. If X, Y, Z be functions of the rectangular co-ordinates x, y, z , dS an element of any limited surface, l, m, n the cosines of the inclinations of the normal at dS to the axes, ds an element of the bounding line, shew that

$$\iint \left\{ l \left(\frac{dZ}{dy} - \frac{dY}{dz} \right) + m \left(\frac{dX}{dz} - \frac{dZ}{dx} \right) + n \left(\frac{dY}{dx} - \frac{dX}{dy} \right) \right\} dS \\ = \int \left(X \frac{dx}{ds} + Y \frac{dy}{ds} + Z \frac{dz}{ds} \right) ds,$$

the differential coefficients of X, Y, Z being partial, and the single integral being taken all round the perimeter of the surface.



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Memorials of Newton and Maxwell

Isaac Newton grave in
Westminster Abbey



https://en.wikipedia.org/wiki/File:Isaac_Newton_grave_in_Westminster_Abbey.jpg

Maxwell's gravestone



https://commons.wikimedia.org/wiki/File:Maxwell%27s_gravestone.JPG

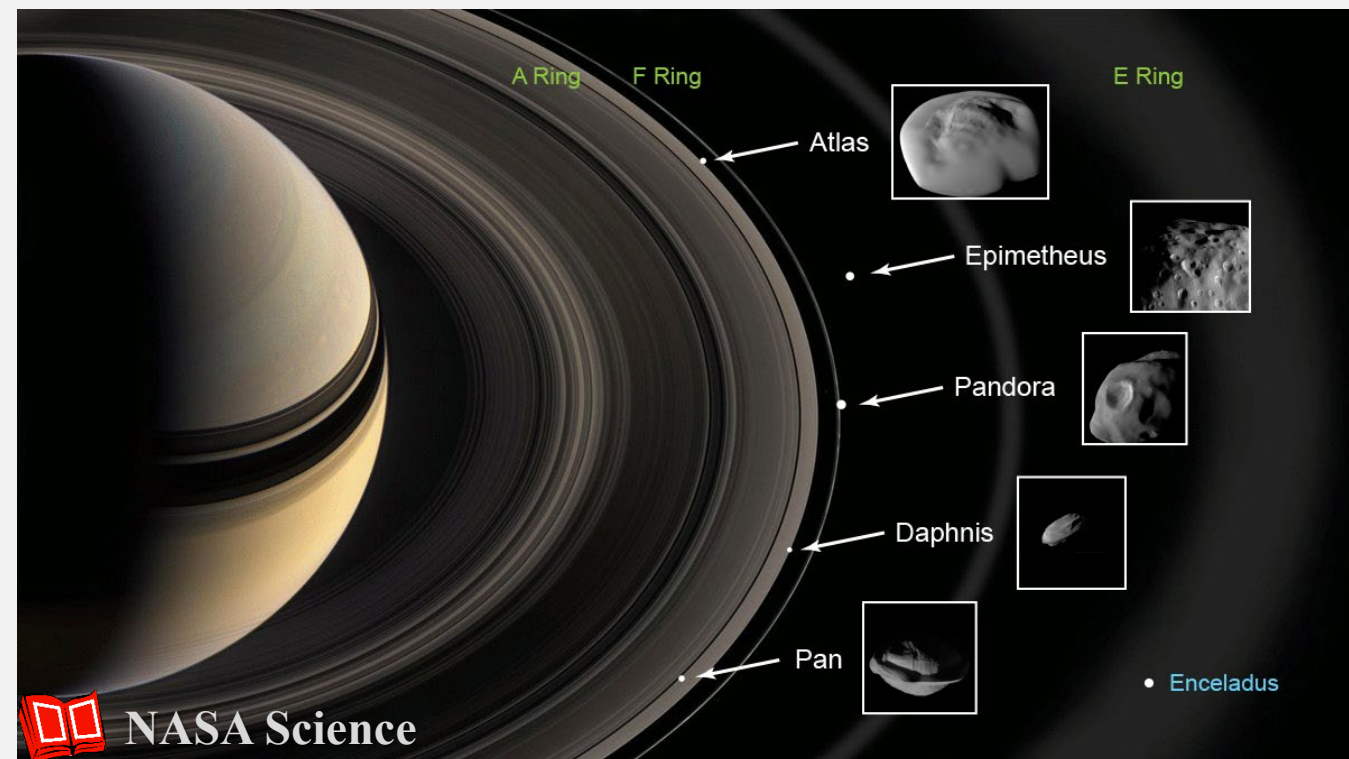
Memorial tablet of
Maxwell in the nave
of Westminster Abbey,
London

James Clerk Maxwell
memorial
This image can be
purchased
from [Westminster Abbey
Library](https://www.westminster-abbey.org/abbey-commemorations/commemorations/james-clerk-maxwell)
Image © 2022 Dean and
Chapter of Westminster

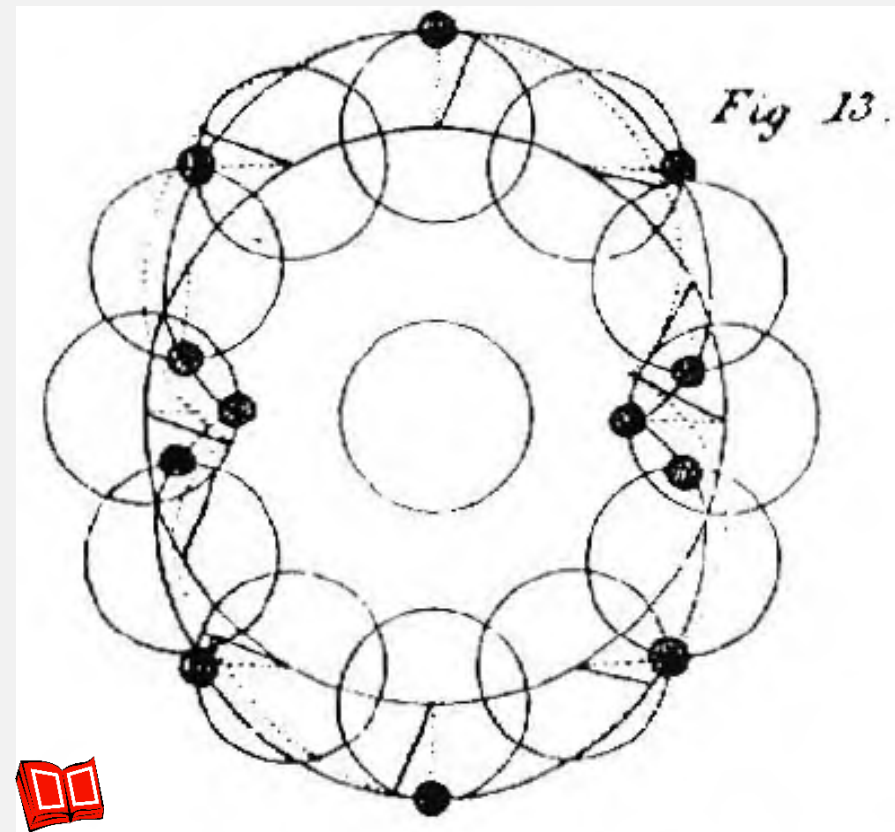
<https://www.westminster-abbey.org/abbey-commemorations/commemorations/james-clerk-maxwell>

Major Research Topics Other Than Electromagnetics

Rings of Saturn



<https://solarsystem.nasa.gov/news/886/cassini-finds-saturns-rings-coat-tiny-moons/>

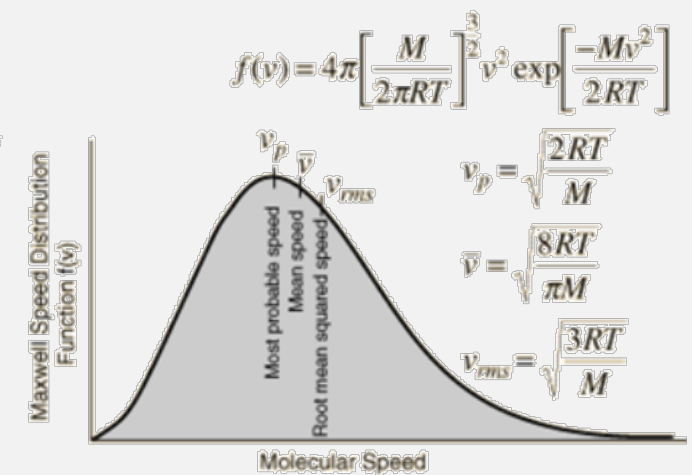
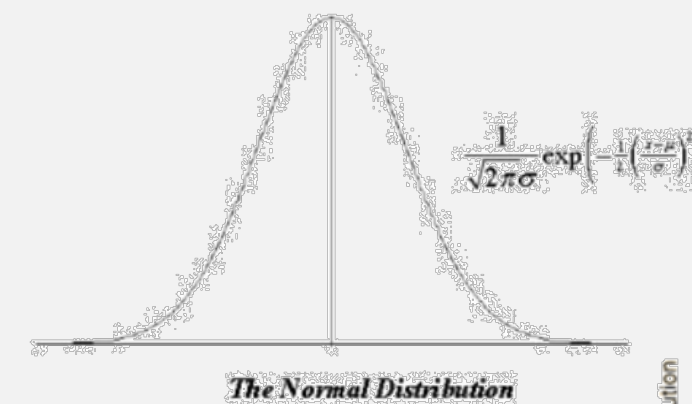


郭奕玲, 沈慧君編, 電磁場理論的奠基人—麥克斯韋, 凡異出版社, 1999, 圖 1-17, p. 3.

Whence we may draw the following conclusions :—

1st. The number of particles whose velocity, resolved in a certain direction, lies between x and $x+dx$ is

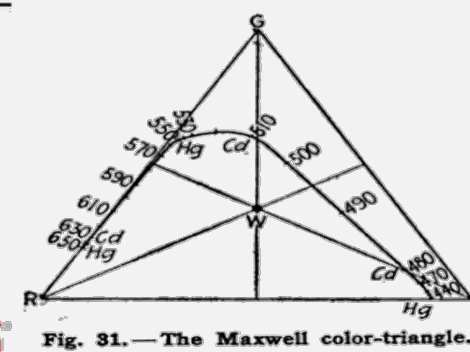
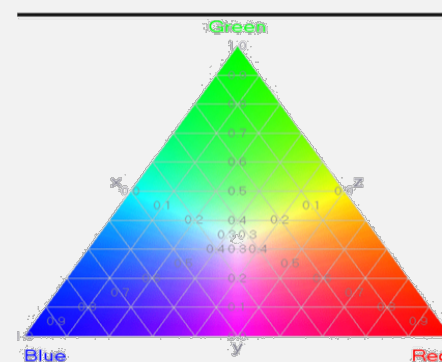
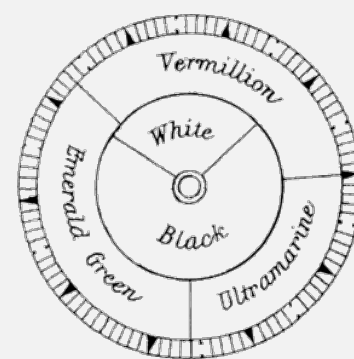
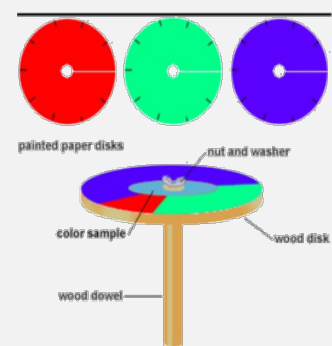
$$N \frac{1}{a\sqrt{\pi}} e^{-\frac{x^2}{a^2}} dx \dots\dots\dots (1).$$



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Colour Vision



SPIE. Digital Library

<https://www.spiedigitallibrary.org/conference-proceedings-of-spice/11099/110990A/Maxwell-color-vision-and-the-color-triangle/10.1117/12.2529364.full?SSO=1>

SPIE. Digital Library

Maxwell and His Equations

A Poem to Tomson's Galvanometer

NATURE

[May 16, 1872]

A LECTURE ON THOMSON'S GALVANO-
METER

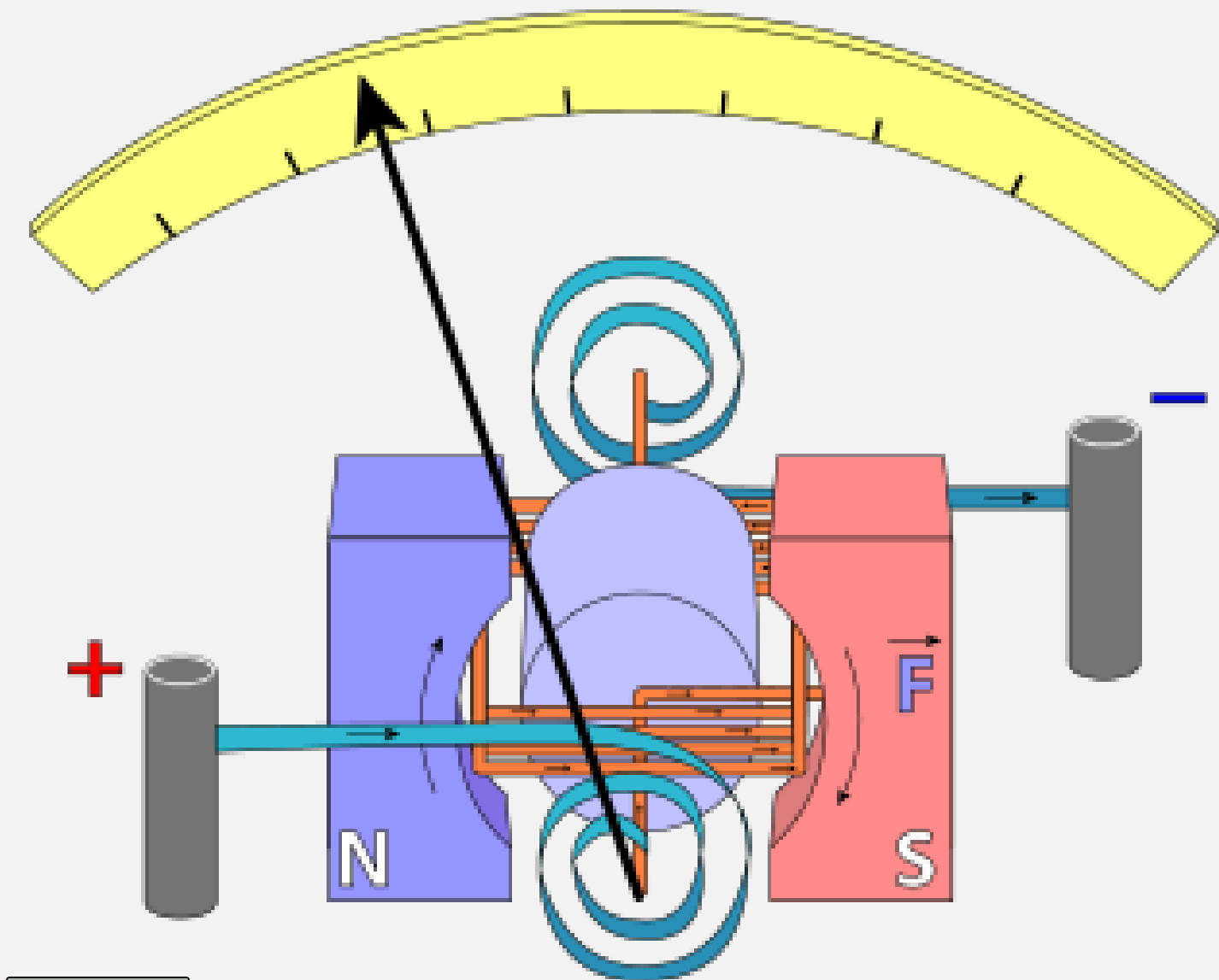
DELIVERED TO A SINGLE PUPIL IN AN ALCOVE WITH
DRAWN CURTAINS

THE lamp-light falls on blackened walls,
And streams through narrow perforations;
The long beam trails o'er pasteboard scales,
With slow-decaying oscillations.
Flow, current! flow! set the quick light-spot flying!
Flow, current! answer, light-spot! flashing, quivering,
dying.

O look! how queer! how thin and clear,
And thinner, clearer, sharper growing,
This gliding fire, with central wire
The fine degrees distinctly showing.
Swing, magnet! swing! advancing and receding;
Swing, magnet! answer, dearest, what's your final reading?

O love! you fail to read the scale
Correct to tenths of a division;
To mirror heaven those eyes were given,
And not for methods of precision.
Break, contact! break! set the free light-spot flying!
Break, contact! rest thee, magnet! swinging, creeping,
dying.

dp/dt  Scottish Science Hall of Fame
<https://digital.nls.uk/scientists/archive/74629736>



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https://commons.wikimedia.org/wiki/File:Galvanometer_scheme.svg

Outline

- Life of James Clark Maxwell
- **Electricity and Magnetism before Maxwell**
- Evolution of Maxwell's Equations
- “Derivation” of Maxwell's Equations

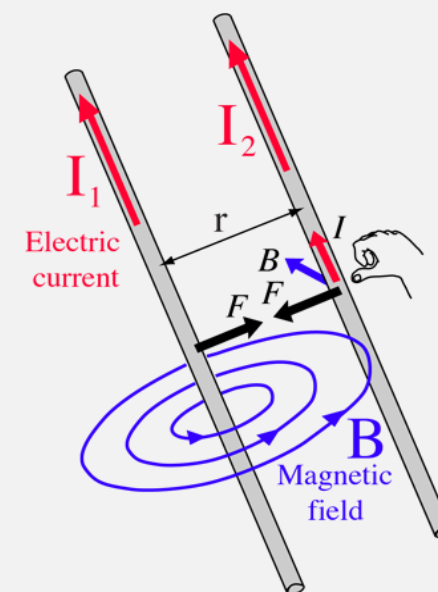
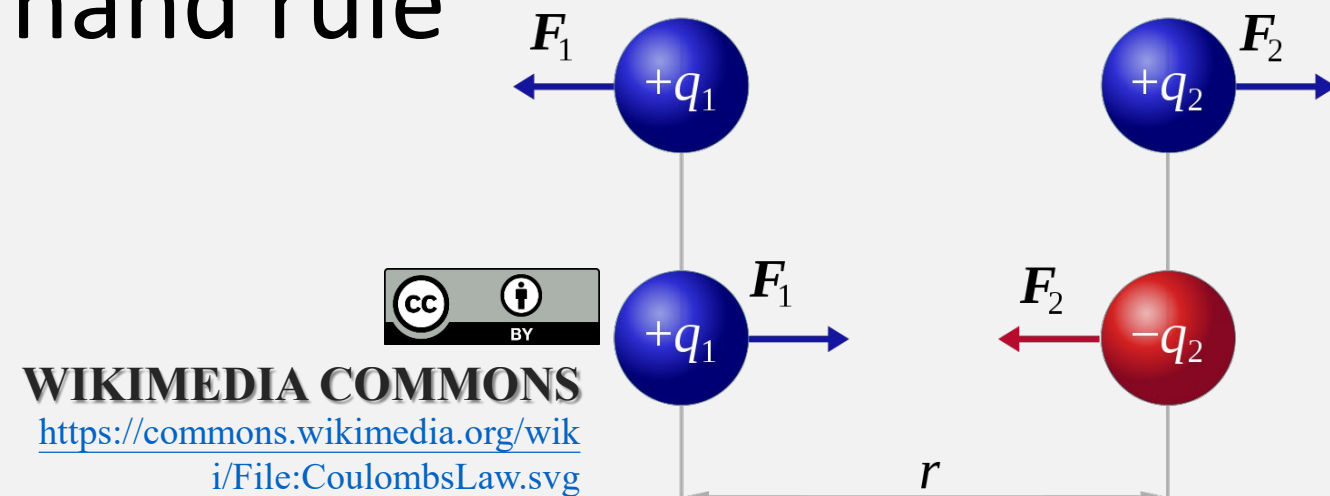
Action at a Distance: Coulomb's Law and Force between Currents

- Coulomb's law for electrostatics

$$F = K \frac{q_1 q_2}{r^2}, \text{ along the direction connecting positions of } q_1 \text{ and } q_2$$

- Force per unit length of parallel currents for magnetostatics

$$F/L = K \frac{I_1 I_2}{r}, \text{ along the direction perpendicular to the direction current and the direction connecting positions of } I_1 \text{ and } I_2 \text{ following the right-hand rule}$$



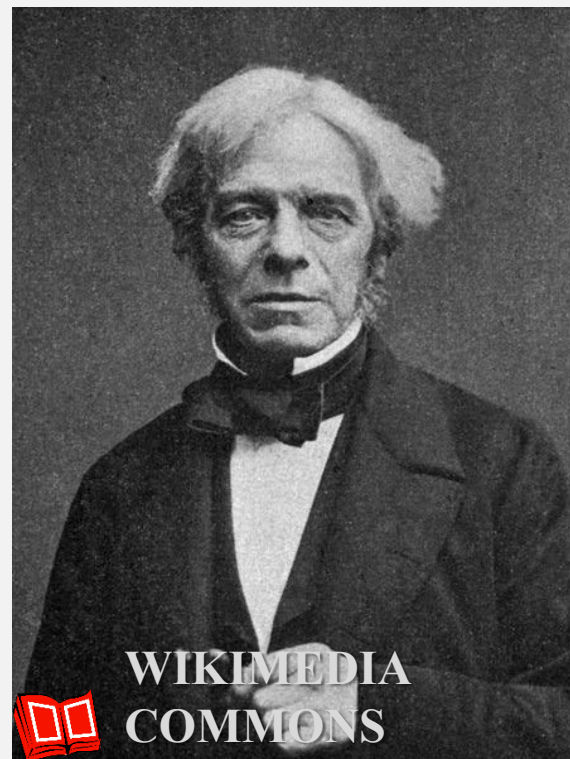
Georgia State University

<http://hyperphysics.phy-astr.gsu.edu/hbase/magnetic/wirfor.html>

Faraday and His Law

- $\varepsilon = - \frac{\Delta \Phi_M}{\Delta t}$

Michael Faraday
1791-1867

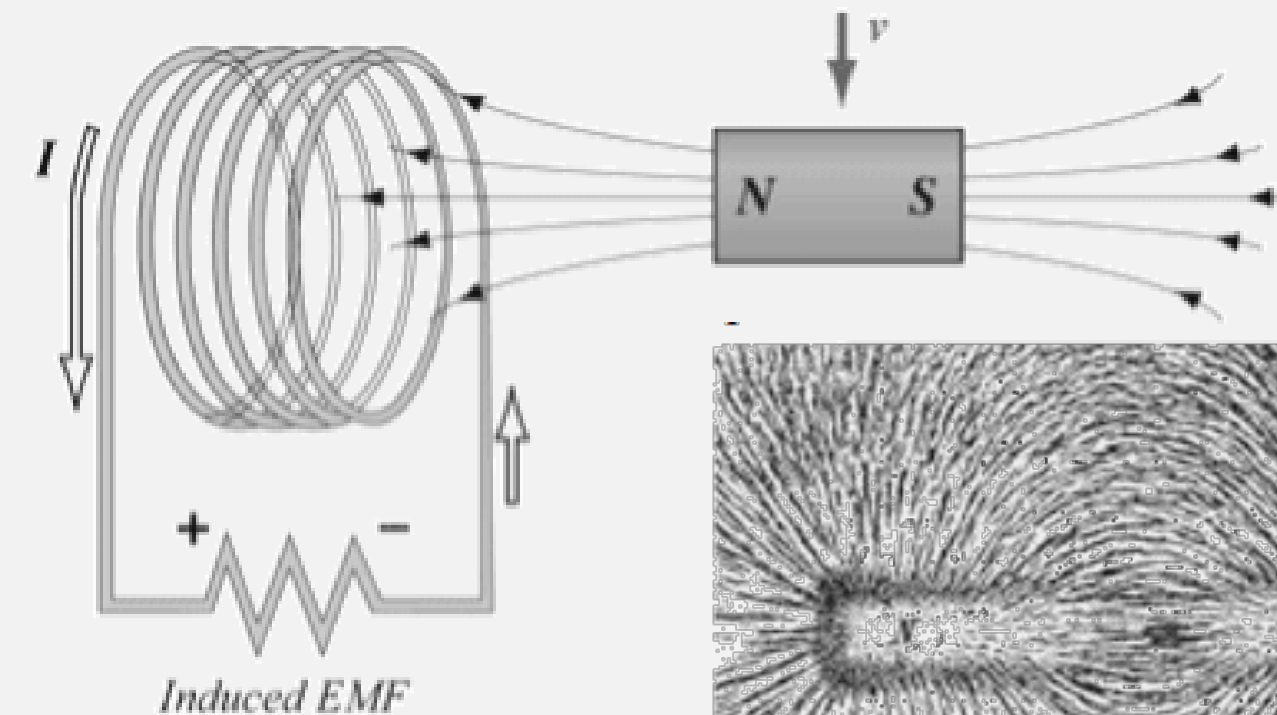


<https://commons.wikimedia.org/wiki/File:Faraday-Millikan-Gale-1913.jpg>

Faraday lecture
at Royal Institution 1856

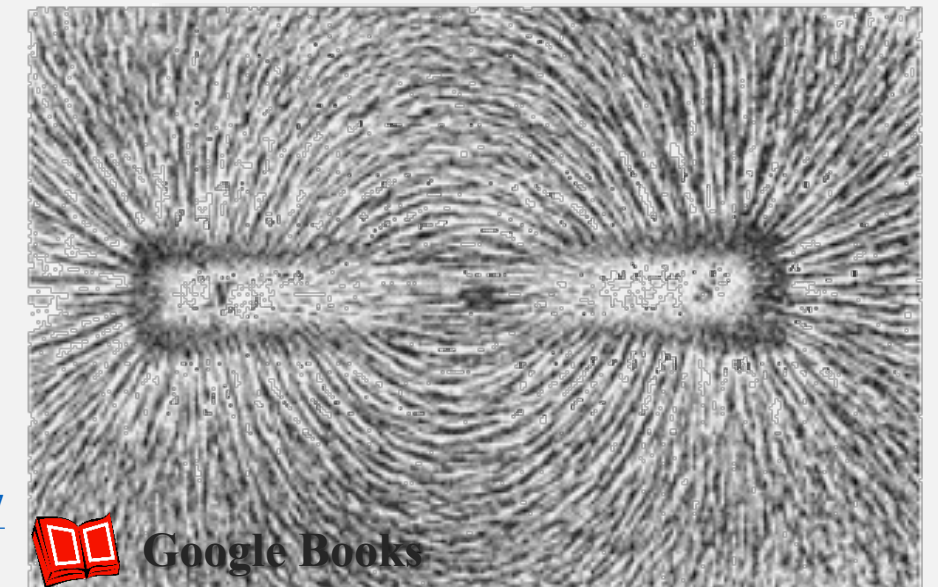


https://commons.wikimedia.org/wiki/File:Faraday_lecture_at_Royal_Institution_1856.jpg



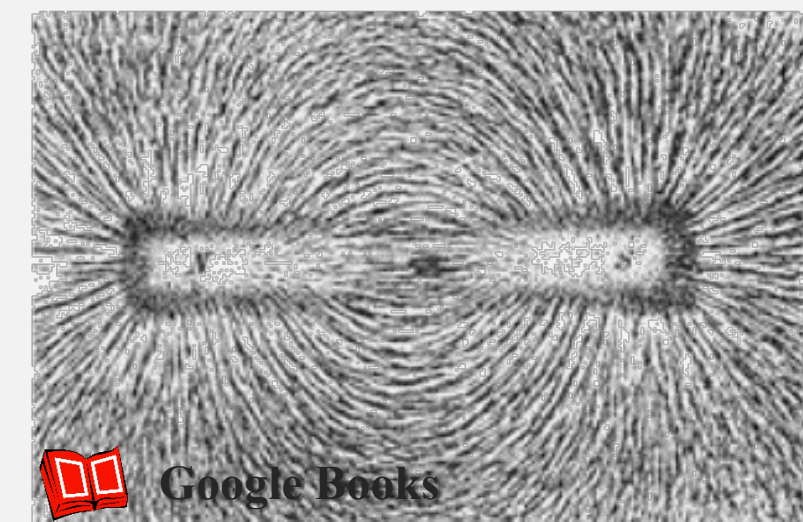
Suffolk University

<https://sites.suffolk.edu/kdshepard/2013/03/07/faradays-law/379a/>



Google Books

https://books.google.com.tw/books?id=BT8AAAAAYA-AJ&pg=PA242&redir_esc=y#v=onepage&q&f=false



https://books.google.com.tw/books?id=BT8A AAAAYAAJ&pg=PA242&redir_esc=y#v=onepage&q&f=false

Faraday's Lines of Force

- ... by line of magnetic force, or **magnetic line of force**, ... I mean that exercise of **magnetic force which is exerted in the lines** usually called magnetic curves, and which equally exist as **passing from or to magnetic poles, or forming concentric circles round an electric current.**
- By **line of electric force**, I mean the force **exerted in the lines joining two bodies**, acting on each other according to the principles of static electric induction (1161, &c), which may also be **either in curved or straight lines.**



Michael Faraday, *Experimental Researches in Electricity*, vol. III, London: Bernard Quaritch, 1855, ¶ 2149.

Outline

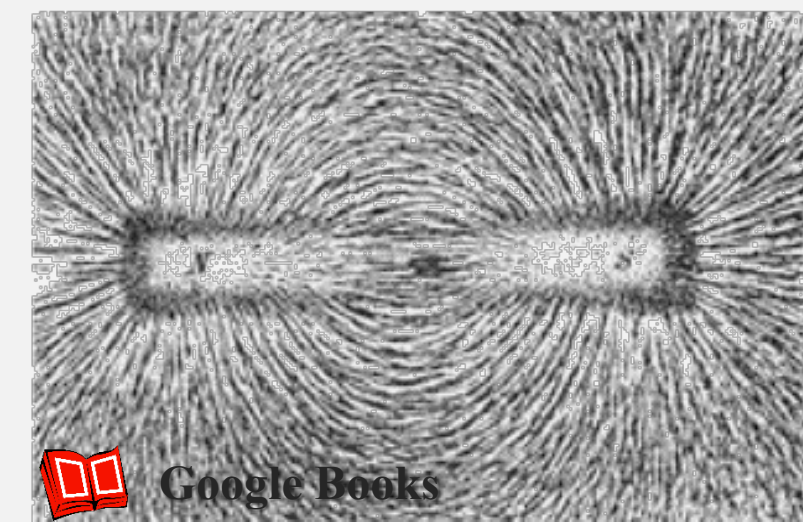
- Life of James Clark Maxwell
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Physical Analogy

- By a physical analogy I mean that **partial similarity** between the laws of one science and those of another which makes each of them illustrate the other
- [A]ll the mathematics sciences are founded on relations between physical laws and laws of numbers, so that the exact science is to **reduce the problems of nature to the determination of quantities by operations on numbers.**



James Clerk Maxwell, On Faraday's Lines of Force, Transactions of the Cambridge Philosophical Society, vol. x, part 1, 1855. p. 155.



Analogy to Incompressible Fluid

- Line of force
 - “[T]his curve will indicate the **direction of that force for every point** through it which it passes, and may called on that account **a line of force** [Maxwell’s emphasis].” (James Clerk Maxwell, On Faraday’s Lines of Force, *Transactions of the Cambridge Philosophical Society*, vol. x, part 1, 1855. p. 158.)

 James Clerk Maxwell, On Faraday’s Lines of Force, *Transactions of the Cambridge Philosophical Society*, vol. x, part 1, 1855. p. 158.

- Tubes of incompressible fluid
 - “[C]onsider these curves not merely lines, but as **fine tubes of variable section** carrying an **incompressible fluid**, . . . We may make the velocity vary according to any given law, by regulating the section of the tube, and in this way we might represent **the intensity of the force as well as its direction by the motion of the fluid in these tubes**.” (James Clerk Maxwell, On Faraday’s Lines of Force, *Transactions of the Cambridge Philosophical Society*, vol. x, part 1, 1855. p. 159.)

 James Clerk Maxwell, On Faraday’s Lines of Force, *Transactions of the Cambridge Philosophical Society*, vol. x, part 1, 1855. p. 159.

Maxwell's Equation, 1865

Equation Names and Notations

Three equations of Magnetic Force	(B)			
„ Electric Currents	(C)			
„ Electromotive Force	(D)			
„ Electric Elasticity	(E)			
„ Electric Resistance	(F)			
„ Total Currents	(A)			
One equation of Free Electricity	(G)			
 „ Continuity	(H)			
	For Electromagnetic Momentum	F	G	H
	„ Magnetic Intensity	α	β	γ
	„ Electromotive Force	P	Q	R
	„ Current due to true conduction	p	q	r
	„ Electric Displacement	f	g	h
	„ Total Current (including variation of displacement) . .	p'	q'	r'
	„ Quantity of free Electricity	e		
	„ Electric Potential	Ψ		

James Clerk Maxwell, “A dynamic theory of the electromagnetic field,” Royal Society Transactions, vol. 155, 1865, pp. 459-512.

Maxwell's Equations (A),(B), (C): 1865

(A)

$$\left. \begin{aligned} p' &= p + \frac{df}{dt}, \\ q' &= q + \frac{dg}{dt}, \\ r' &= r + \frac{dh}{dt}, \end{aligned} \right\}$$


(B)

$$\left. \begin{aligned} \mu\alpha &= \frac{dH}{dy} - \frac{dG}{dz}, \\ \mu\dot{\beta} &= \frac{dF}{dz} - \frac{dH}{dx}, \\ \mu\gamma &= \frac{dG}{dx} - \frac{dF}{dy}. \end{aligned} \right\}$$


(C)

$$\left. \begin{aligned} \frac{d\gamma}{dy} - \frac{d\beta}{dz} &= 4\pi p', \\ \frac{d\alpha}{dz} - \frac{d\gamma}{dx} &= 4\pi q', \\ \frac{d\beta}{dx} - \frac{d\alpha}{dy} &= 4\pi r'. \end{aligned} \right\}$$


James Clerk Maxwell, "A dynamic theory of the electromagnetic field," Royal Society Transactions, vol. 155, 1865, pp. 459-512.

Maxwell's Equations (D),(E), (F), (G), (H): 1865

(D)

$$\left. \begin{aligned} P &= \mu \left(\gamma \frac{dy}{dt} - \beta \frac{dz}{dt} \right) - \frac{dF}{dt} - \frac{d\Psi}{dx}, \\ Q &= \mu \left(\alpha \frac{dz}{dt} - \gamma \frac{dx}{dt} \right) - \frac{dG}{dt} - \frac{d\Psi}{dy}, \\ R &= \mu \left(\beta \frac{dx}{dt} - \alpha \frac{dy}{dt} \right) - \frac{dH}{dt} - \frac{d\Psi}{dz}. \end{aligned} \right\}$$

(E)

$$\left. \begin{aligned} P &= kf, \\ Q &= kg, \\ R &= kh. \end{aligned} \right\}$$

(F)

$$\left. \begin{aligned} P &= -\xi p, \\ Q &= -\xi q, \\ R &= -\xi r. \end{aligned} \right\}$$

(G)

$$e + \frac{df}{dx} + \frac{dg}{dy} + \frac{dh}{dz} = 0.$$

(H)

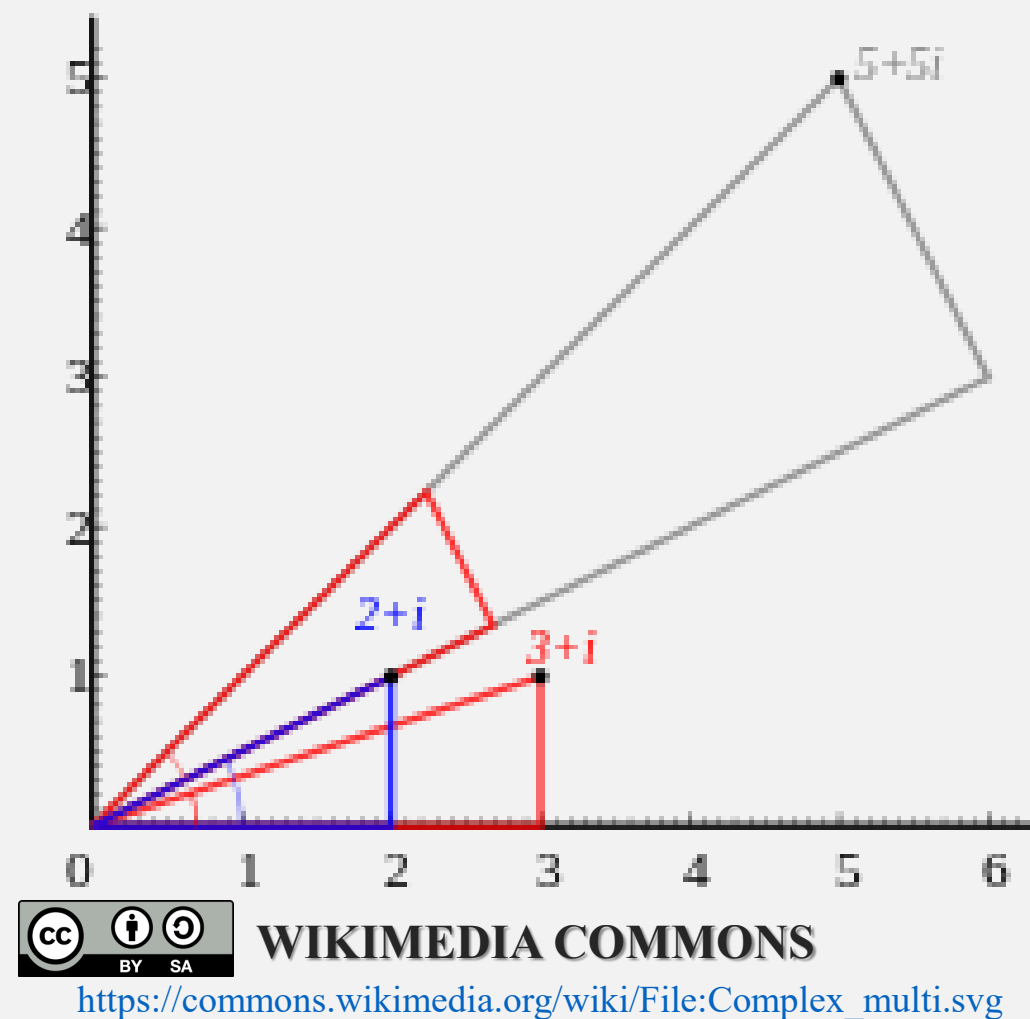
$$\frac{de}{dt} + \frac{dp}{dx} + \frac{dq}{dy} + \frac{dr}{dz} = 0$$

James Clerk Maxwell, "A dynamic theory of the electromagnetic field," Royal Society Transactions, vol. 155, 1865, pp. 459-512.

Complex Number and Quaternion

$$q = s + x \mathbf{i} + y \mathbf{j} + z \mathbf{k}$$

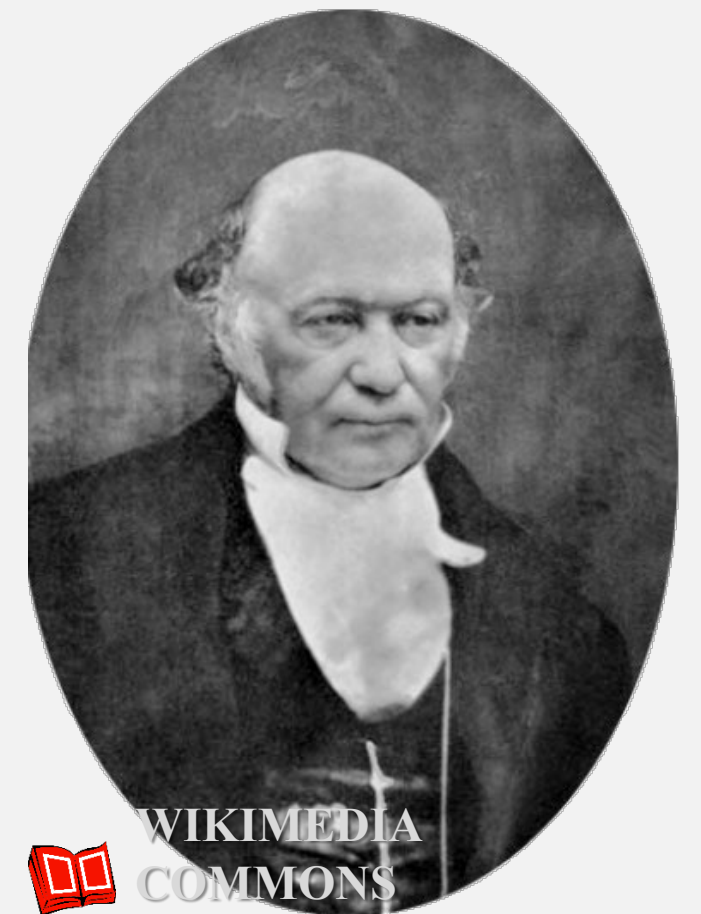
Special case: $s = 0$



$\times$	1	i	j	k
1	1	i	j	k
i	i	-1	k	$-j$
j	j	$-k$	-1	i
k	k	j	$-i$	-1

WIKIPEDIA
<https://en.wikipedia.org/wiki/Quaternion>

$$q_1 q_2 = -(x_1 x_2 + y_1 y_2 + z_1 z_2) + (y_1 z_2 - z_1 y_2) \mathbf{i} + (z_1 x_2 - x_1 z_2) \mathbf{j} + (x_1 y_2 - x_2 y_1) \mathbf{k}$$



Sir William Rowan Hamilton (1805–1865)

[https://commons.wikimedia.org/wiki/File:William Rowan Hamilton portrait oval combined.png](https://commons.wikimedia.org/wiki/File:William_Rowan_Hamilton_portrait_oval_combined.png)

Maxwell's Equations, 1873 and Corresponding Modern Forms, (A),(B),(C)

$$\begin{aligned}
 a &= \frac{dH}{dy} - \frac{dG}{dz} \\
 b &= \frac{dF}{dz} - \frac{dH}{dx} \\
 c &= \frac{dG}{dx} - \frac{dF}{dy} \\
 P &= c \frac{dy}{dt} - b \frac{dz}{dt} - \frac{dF}{dt} - \frac{d\psi}{dx} \\
 Q &= a \frac{dz}{dt} - c \frac{dx}{dt} - \frac{dG}{dt} - \frac{d\psi}{dy} \\
 R &= b \frac{dx}{dt} - a \frac{dy}{dt} - \frac{dH}{dt} - \frac{d\psi}{dz} \\
 X &= vc - wb \\
 Y &= wa - uc \\
 Z &= ub - va
 \end{aligned}
 \tag{A} \qquad \mathbf{B} = \nabla \times \mathbf{A}$$

$$\begin{aligned}
 P &= c \frac{dy}{dt} - b \frac{dz}{dt} - \frac{dF}{dt} - \frac{d\psi}{dx} \\
 Q &= a \frac{dz}{dt} - c \frac{dx}{dt} - \frac{dG}{dt} - \frac{d\psi}{dy} \\
 R &= b \frac{dx}{dt} - a \frac{dy}{dt} - \frac{dH}{dt} - \frac{d\psi}{dz} \\
 X &= vc - wb \\
 Y &= wa - uc \\
 Z &= ub - va
 \end{aligned}
 \tag{B} \qquad \mathbf{E} = \mathbf{v} \times \mathbf{B} - \frac{\partial \mathbf{A}}{\partial t} - \nabla \phi$$

$$\begin{aligned}
 X &= vc - wb \\
 Y &= wa - uc \\
 Z &= ub - va
 \end{aligned}
 \tag{C} \qquad \mathbf{F} = \mathbf{J} \times \mathbf{B}$$



Tapan K. Sarkar, etc., History of Wireless, New Jersey: John Wiley & Sons, Inc., 2006, Fig. 3.11, p. 185.

James Clerk Maxwell, A Treatise on Electricity and Magnetism, Oxford: Clarendon Press, 1873, pp. 480-493.

Maxwell's Equations, 1873 and Corresponding Modern Forms, (D),(E),(F),(G)

$$a = \alpha + 4\pi A$$

$$b = \beta + 4\pi B$$

$$c = \gamma + 4\pi C$$

$$4\pi u = \frac{d\gamma}{dy} - \frac{d\beta}{dz}$$

$$4\pi v = \frac{d\alpha}{dz} - \frac{d\gamma}{dx}$$

$$4\pi w = \frac{d\beta}{dx} - \frac{d\alpha}{dy}$$

$$\mathfrak{D} = \frac{1}{4\pi} K \mathfrak{E}$$

$$\mathfrak{H} = C \mathfrak{E}$$

(D)

$$\mathbf{B} = \mu_0 \mathbf{H} + \mathbf{M}$$

(E)

$$\mathbf{J} = \nabla \times \mathbf{H}$$

(F)

$$\mathbf{D} = \epsilon \mathbf{E}$$

(G)

$$\mathbf{J}_c = \sigma \mathbf{E}$$



Tapan K. Sarkar, etc., History of Wireless, New Jersey: John Wiley & Sons, Inc., 2006, Fig. 3.11, p. 185.

James Clerk Maxwell, A Treatise on Electricity and Magnetism, Oxford: Clarendon Press, 1873, pp. 480-493.

Maxwell's Equations, 1873 and Corresponding Modern Forms, (H), (H*), (I), (I*), (J), (K), (L)

$$\mathcal{C} = (C + \frac{1}{4\pi} K \frac{d}{dt}) \mathfrak{C}$$

(I)

$$\mathbf{J} = \sigma \mathbf{E} + \epsilon \frac{\partial \mathbf{E}}{\partial t}$$

$$u = CP + \frac{1}{4\pi} K \frac{dP}{dt}$$

$$v = CQ + \frac{1}{4\pi} K \frac{dQ}{dt}$$

$$w = CR + \frac{1}{4\pi} K \frac{dR}{dt}$$

(I*)

$$\mathbf{J} = \sigma \mathbf{E} + \epsilon \frac{\partial \mathbf{E}}{\partial t}$$

$$\rho = \frac{df}{dx} + \frac{dg}{dy} + \frac{dh}{dz}$$

(J)

$$\varrho = \nabla \cdot \mathbf{D}$$

$$\sigma = lf + mg + nh + l'f' + m'g' + n'h'$$

(K)

$$\varrho_s = \mathbf{n} \cdot (\mathbf{D}_1 - \mathbf{D}_2)$$

$$\mathfrak{B} = \mu \mathfrak{H}$$

(L)

$$\mathbf{B} = \mu \mathbf{H}$$



Tapan K. Sarkar, etc., History of Wireless, New Jersey: John Wiley & Sons, Inc., 2006, Fig. 3.11, p. 185.
James Clerk Maxwell, A Treatise on Electricity and Magnetism, Oxford: Clarendon Press, 1873, pp. 480-493.

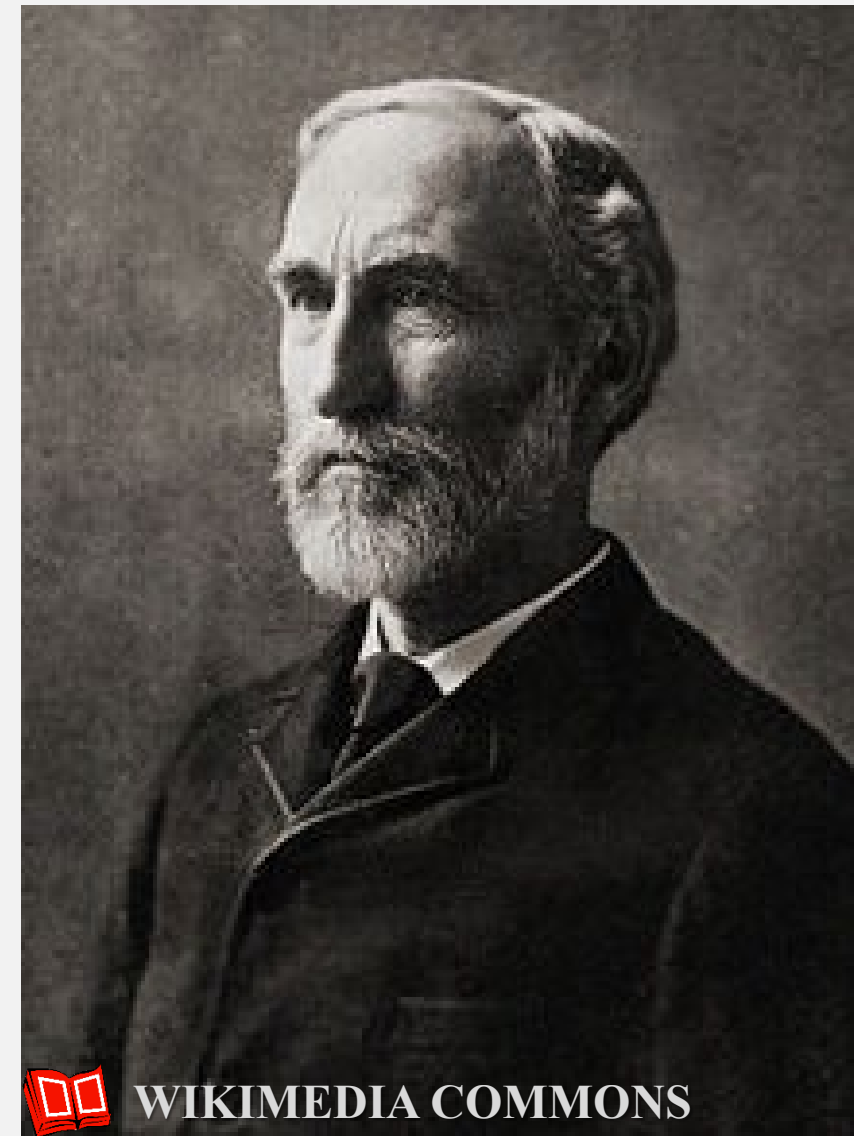
Vector Analysis: Oliver Heaviside and Josiah Willard Gibbs



 WIKIMEDIA COMMONS

Oliver Heaviside
1850–1925

<https://commons.wikimedia.org/wiki/File:Oheaviside.jpg>



 WIKIMEDIA COMMONS

Josiah Willard Gibbs
1839–1903

https://commons.wikimedia.org/wiki/File:Josiah_Willard_Gibbs_-from_MMS-.jpg

- Scalar product (Oliver Heaviside, *Electromagnetic Theory*, New York: Dover Publications, Inc., 1950, p. 38, ¶ 107.)

 • $\mathbf{AB} = AB \cos \theta$
Oliver Heaviside, *Electromagnetic Theory*, New York: Dover Publications, Inc., 1950, p. 38, ¶ 107.

- Vector product (Oliver Heaviside, *Electromagnetic Theory*, New York: Dover Publications, Inc., 1950, p. 40, ¶ 111.)

 • $\mathbf{C} = V\mathbf{AB}, C = AB \sin \theta$
Oliver Heaviside, *Electromagnetic Theory*, New York: Dover Publications, Inc., 1950, p. 40, ¶ 111.

Oliver Heaviside's Circuitation, curl, and divergence, 1885, 1893

- Circuitation (Oliver Heaviside, *Electromagnetic Theory*, New York: Dover Publications, Inc., 1950, p. 9, ¶ 33.)
 - “used,..., to indicate the often-occurring operation of a line-integral in a closed circuit”



Oliver Heaviside, *Electromagnetic Theory*, New York: Dover Publications, Inc., 1950, p. 9, ¶ 33.

- curl (Oliver Heaviside, *Electromagnetic Theory*, New York: Dover Publications, Inc., 1950, p. 9, ¶ 36.)
 - “The circuitation of the force is then called its ‘curl’.”



Oliver Heaviside, *Electromagnetic Theory*, New York: Dover Publications, Inc., 1950, p. 9, ¶ 36.

- divergence (Oliver Heaviside, *Electromagnetic Theory*, New York: Dover Publications, Inc., 1950, p. 13, ¶ 51.)
 - “The divergence of any flux is the amount of the flux leaving the unit volume.”



Oliver Heaviside, *Electromagnetic Theory*, New York: Dover Publications, Inc., 1950, p. 13, ¶ 51.

Oliver Heaviside's Duplex Form of Maxwell's Equations, 1885, 1893

- $\mathbf{D} = \epsilon \mathbf{E}$
- $\mathbf{B} = \mu \mathbf{H}$
- $\text{curl } \mathbf{H} = \mathbf{J}$
- $-\text{curl } \mathbf{E} = \mathbf{G}$
- $\text{div } \mathbf{D} = \rho$
- $\text{div } \mathbf{B} = \sigma$

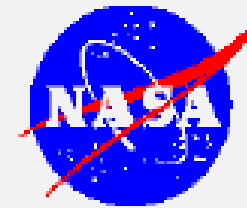
Maxwell's Equations, Modern Text Version

- $\nabla \cdot \bar{D}(\bar{r}, t) = \rho_e(\bar{r}, t)$
- $\nabla \cdot \bar{B}(\bar{r}, t) = 0$
- $\nabla \times \bar{E}(\bar{r}, t) = -\frac{\partial}{\partial t} \bar{B}(\bar{r}, t)$
- $\nabla \times \bar{H}(\bar{r}, t) = \bar{J}(\bar{r}, t) + \frac{\partial}{\partial t} \bar{D}(\bar{r}, t)$

Outline

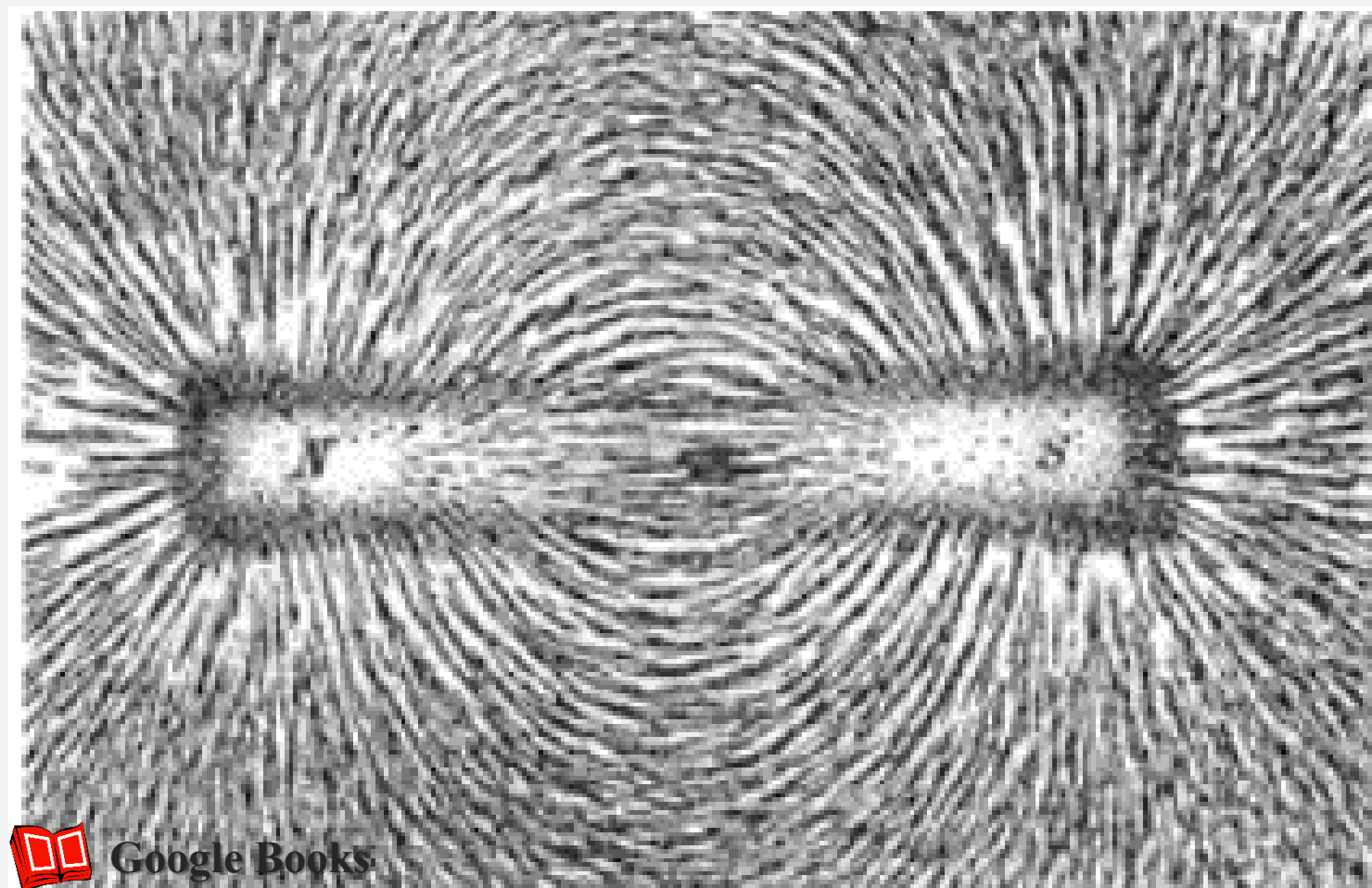
- Life of James Clark Maxwell
- Electricity and Magnetism before Maxwell
- Evolution of Maxwell's Equations
- **“Derivation” of Maxwell's Equations**

Lines of Force and Stream Lines



Definition of Streamlines

Glenn
Research
Center



Google Books

https://books.google.com.tw/books?id=BT8AAAAAYAAJ&pg=PA242&redir_esc=y#v=onepage&q&f=false



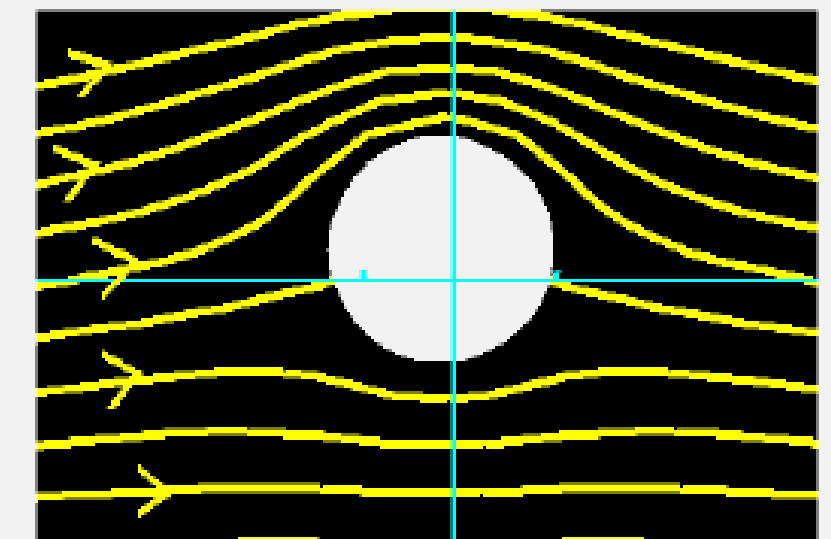
Airfoil

Moving with the object.
Flow goes from left to right.
Steady, two-dimensional flow.

Streamline is a path traced out by a massless particle moving with the flow.

Velocity is tangent to streamline at every point.

Mass does not cross streamlines.

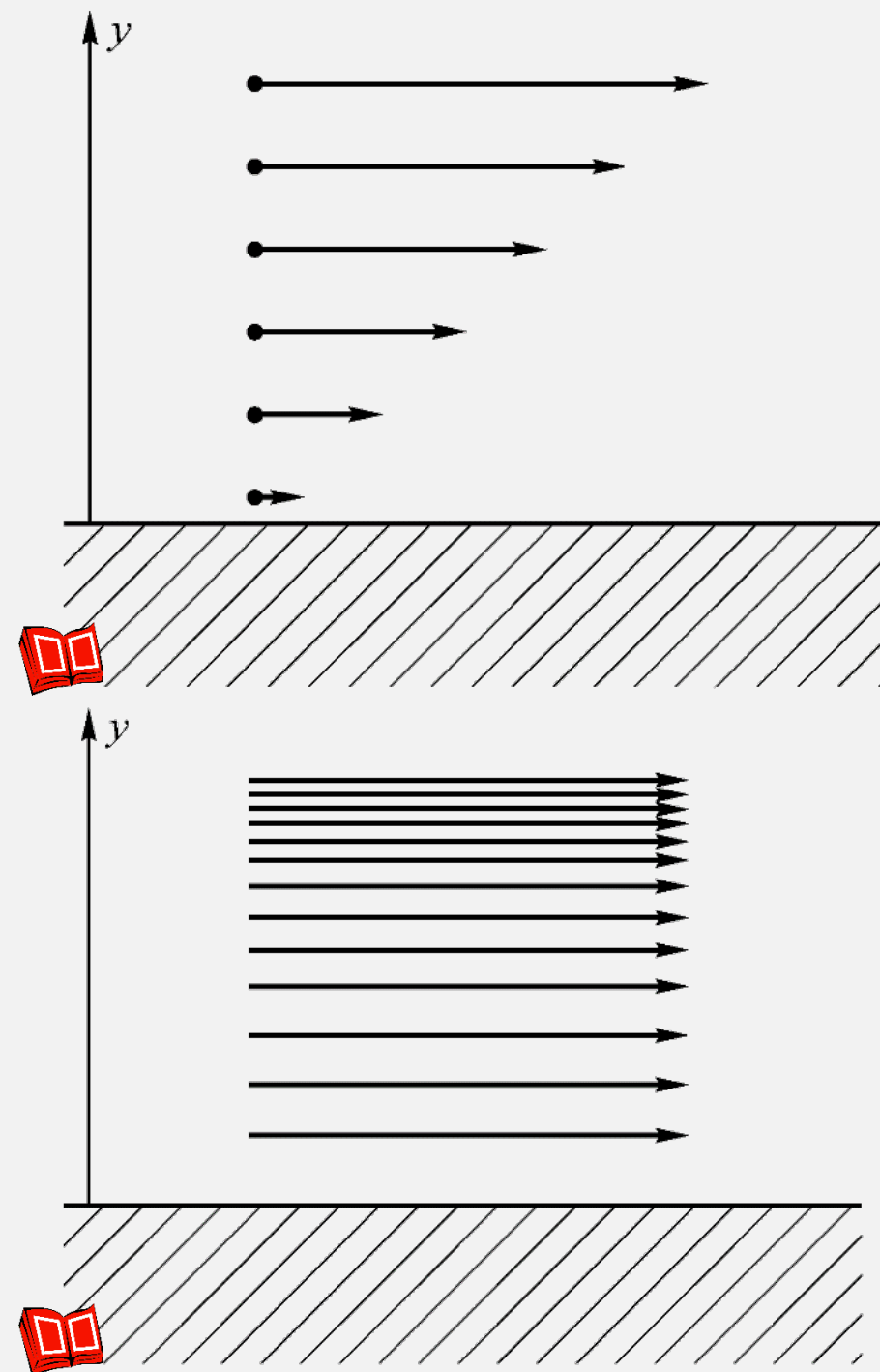


Cylinder

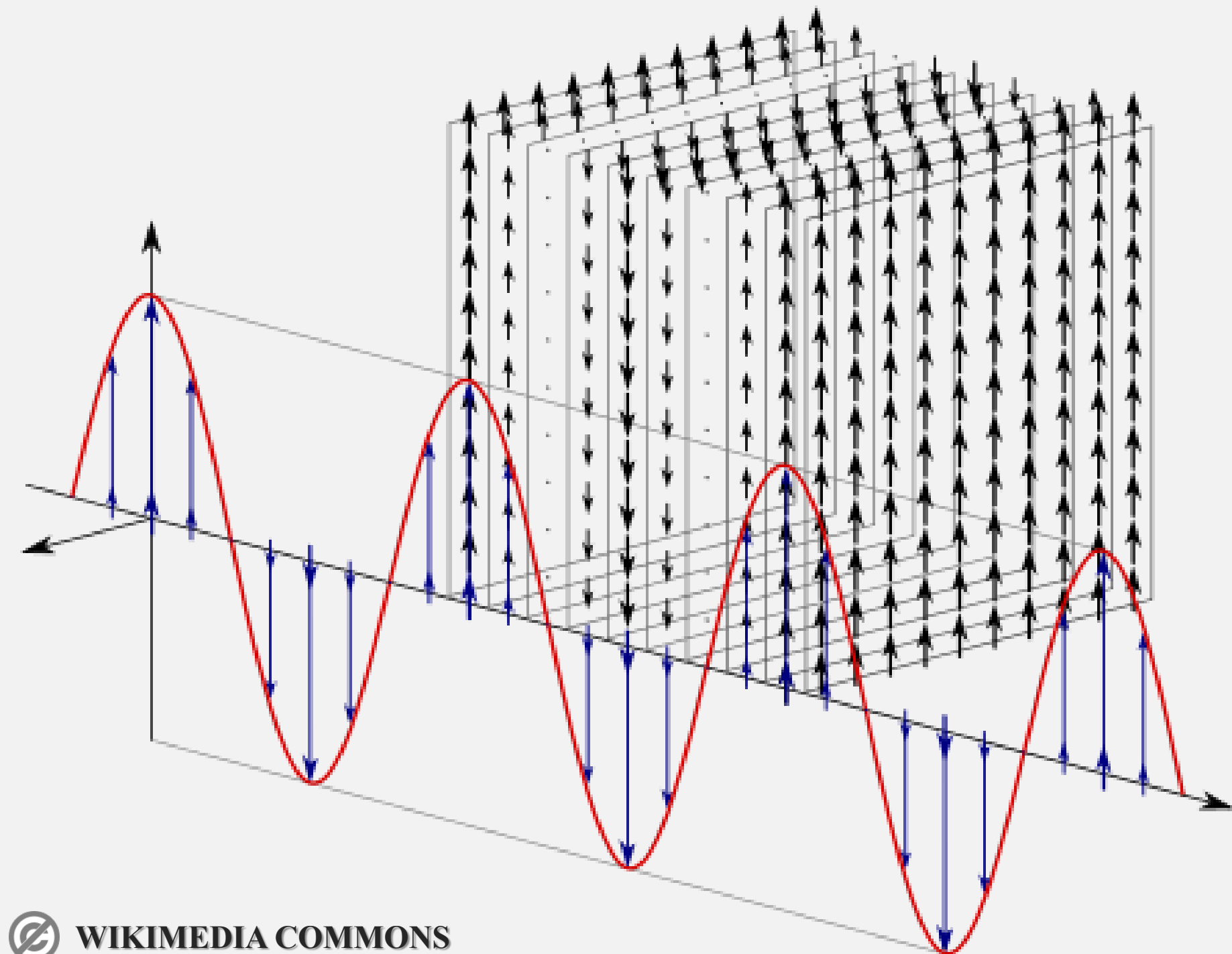


NASA <https://www.grc.nasa.gov/WWW/k-12/VirtualAero/BottleRocket/airplane/stream.html>

Expressions of Field Strength with Lines of Field



鄭士康, 電磁波, 台北市: 全華, 2006, 附錄 B.

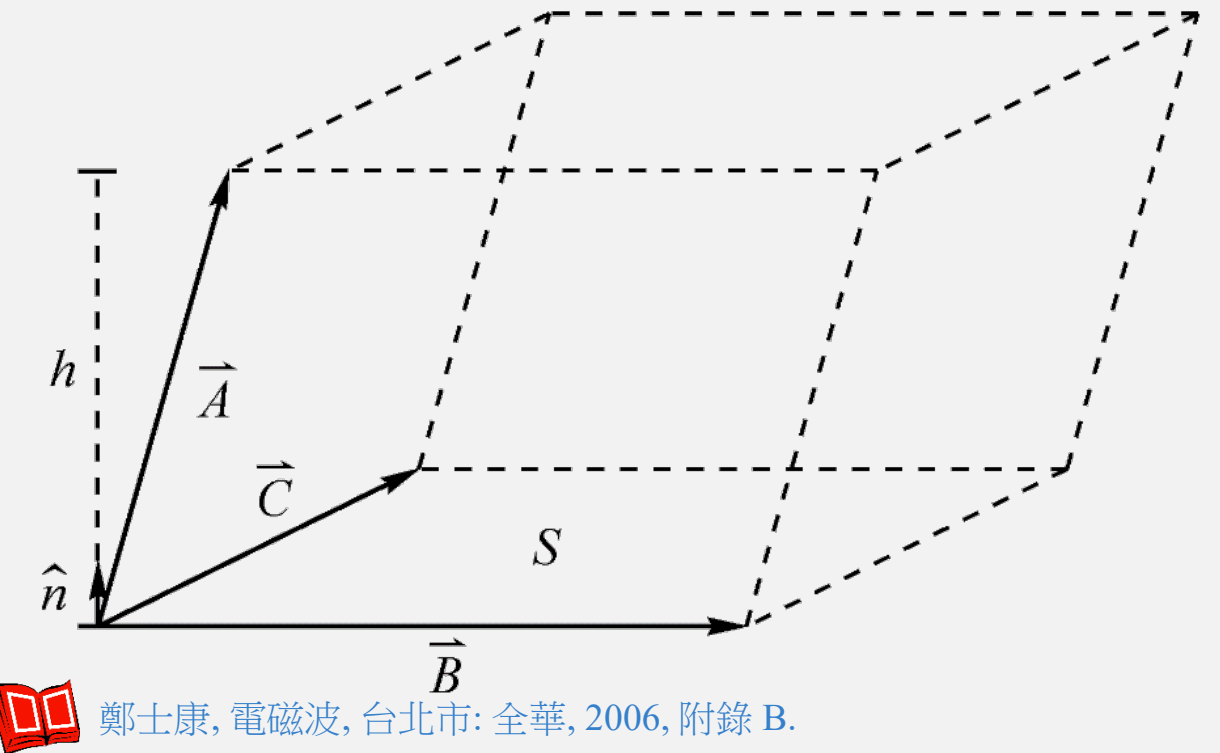


WIKIMEDIA COMMONS

https://commons.wikimedia.org/wiki/File:Linear_Polarization_Linearly_Polarized_Light_plane_wave.svg

Vector, Products, and Magnitude

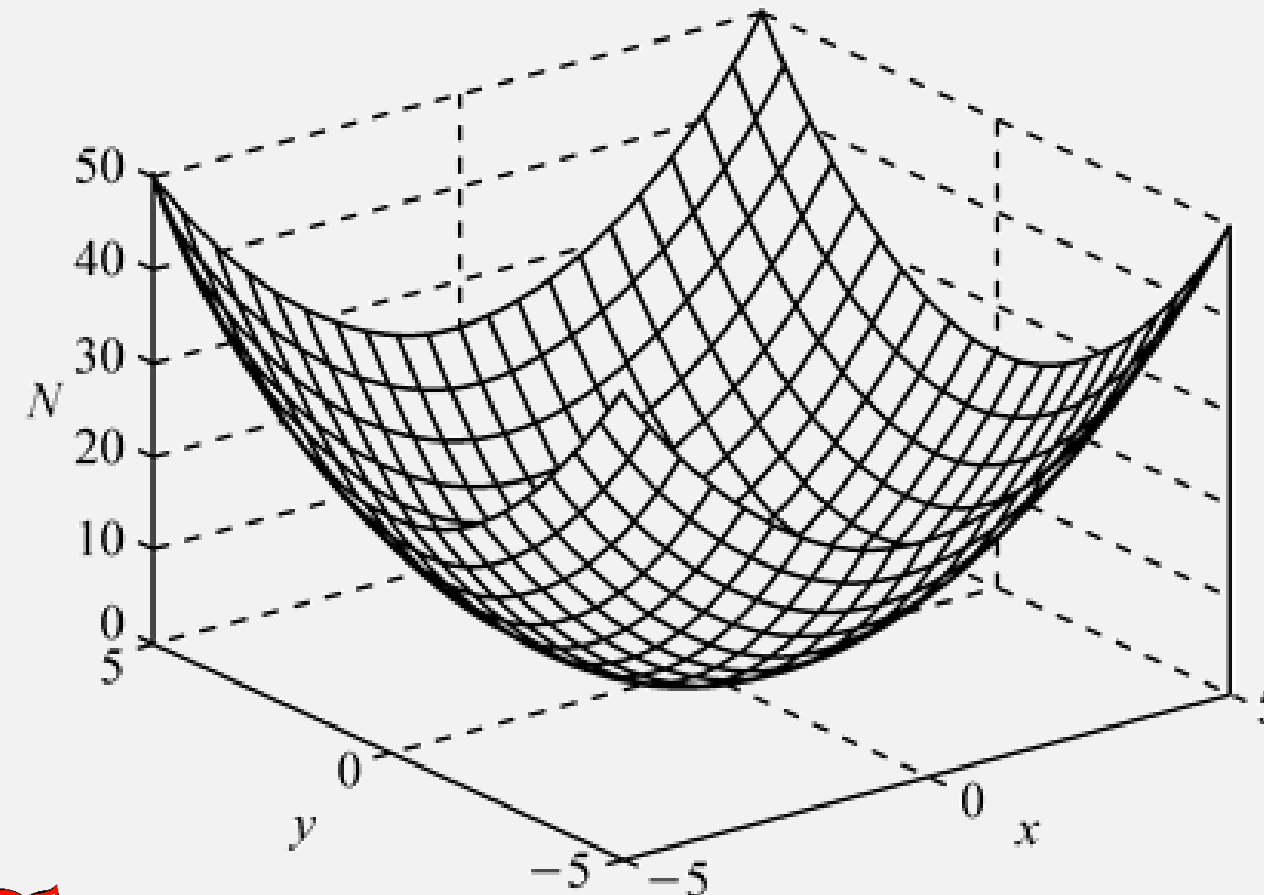
- $\vec{V} = [V_x \quad V_y \quad V_z]^T = \hat{x}V_x + \hat{y}V_y + \hat{z}V_z$
- $|\vec{V}| = \sqrt{V_x^2 + V_y^2 + V_z^2}$
- $\vec{A} \cdot \vec{B} = A_x B_x + A_y B_y + A_z B_z = |\vec{A}| |\vec{B}| \cos \theta_{AB}$
- $|\vec{V}| = \sqrt{\vec{V} \cdot \vec{V}}$
- $\vec{B} \times \vec{C} = \hat{x}(B_y C_z - B_z C_y) + \hat{y}(B_z C_x - B_x C_z) + \hat{z}(B_x C_y - B_y C_x)$
- $\vec{B} \times \vec{C} = \begin{vmatrix} \hat{x} & \hat{y} & \hat{z} \\ B_x & B_y & B_z \\ C_x & C_y & C_z \end{vmatrix} = \hat{n} |\vec{B}| |\vec{C}| \sin \theta_{BC}$



鄭士康, 電磁波, 台北市: 全華, 2006, 附錄 B.

Position Vector and Surface

- $\bar{r} = \hat{x}x + \hat{y}y + \hat{z}z$
- $\bar{r} = \bar{r}(u, v) = \hat{x}x(u, v) + \hat{y}y(u, v) + \hat{z}z(u, v)$
- Example: $\bar{r}(u, v) = \hat{x}u + \hat{y}v + \hat{z}(u^2 + v^2)$
- $\hat{n}(x, y, z) = \hat{n}(u, v) =$
- $\hat{x}n_x + \hat{y}n_y + \hat{z}n_z =$
- $\frac{\partial \bar{r}(u, v)}{\partial u} \times \frac{\partial \bar{r}(u, v)}{\partial v} / \text{Jacobian}$
- $\text{Jacobian} = \left| \frac{\partial \bar{r}(u, v)}{\partial u} \times \frac{\partial \bar{r}(u, v)}{\partial v} \right|$



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4

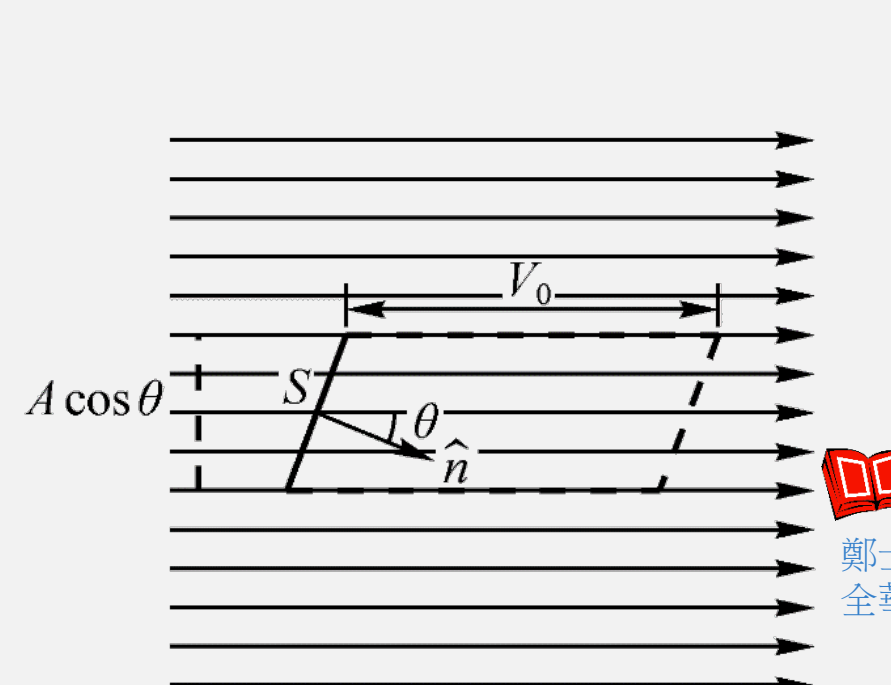
Flux and Flux Density

- Flux density

$$\bar{V}(x, y, z, t) = \hat{x}V_x(x, y, z, t) + \hat{y}V_y(x, y, z, t) + \hat{z}V_z(x, y, z, t)$$

- Flux through surface S

$$\Phi(t) = \int_S \bar{V}(x, y, z, t) \cdot \hat{n} da$$



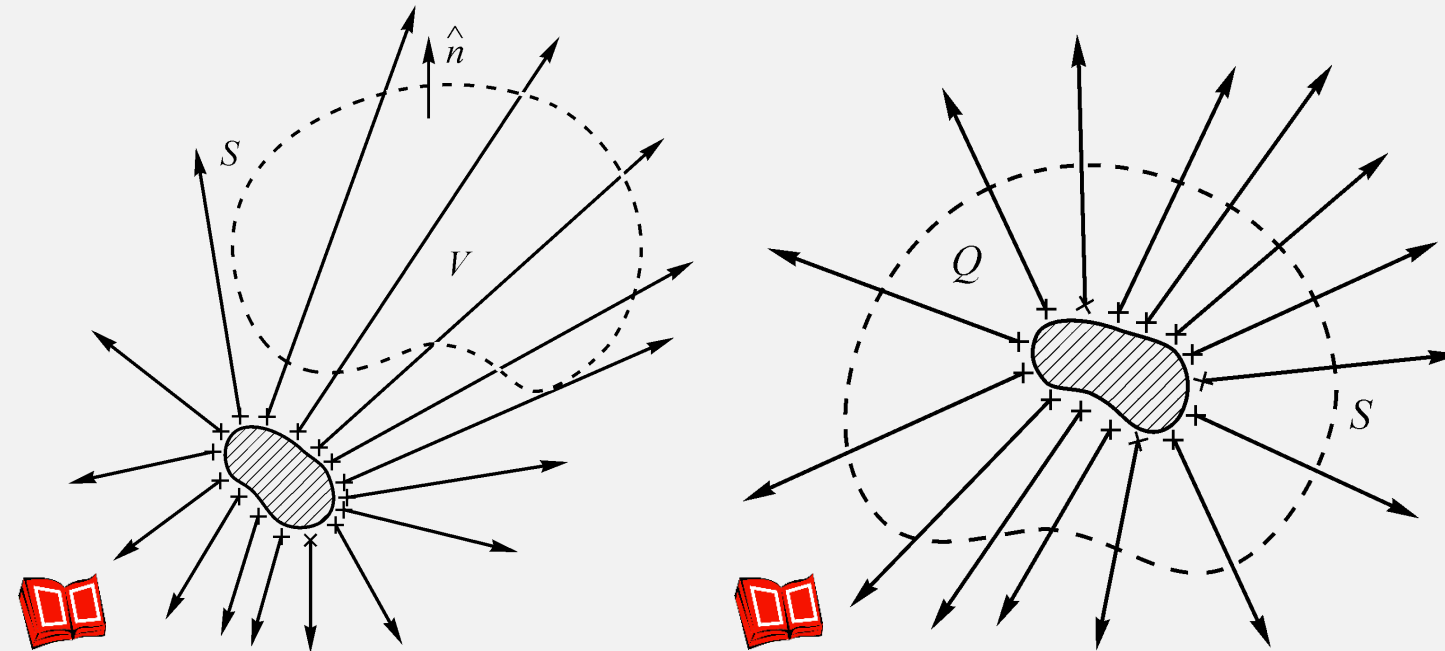
鄭士康, 電磁波, 台北市:
全華, 2006, 附錄 B.

- $\Phi(t) =$

$$\iint_{S_{yz}} V_x(x, y, z, t) n_x dy dz + \iint_{S_{zx}} V_y(x, y, z, t) n_y dz dx + \iint_{S_{xy}} V_z(x, y, z, t) n_z dx dy$$

Electric Flux Density and Gauss's Law

- Lines of electric field
- Imaginary substance in the analogy to an imaginary incompressible fluid
- Electric flux Φ_E and electric charge
- Flux tube
- Electric Flux Density $\bar{D}(x, y, z, t)$
- Gauss's law



鄭士康, 電磁波, 台北市: 全華, 2006, 附錄 B.

$$Q(t) = \oint_S \bar{D}(x, y, z, t) \cdot \hat{n} da = \iint_{S_{yz}} D_x(x, y, z, t) n_x dy dz + \iint_{S_{zx}} D_y(x, y, z, t) n_y dz dx + \iint_{S_{xy}} D_z(x, y, z, t) n_z dx dy$$

Divergence (1/2)

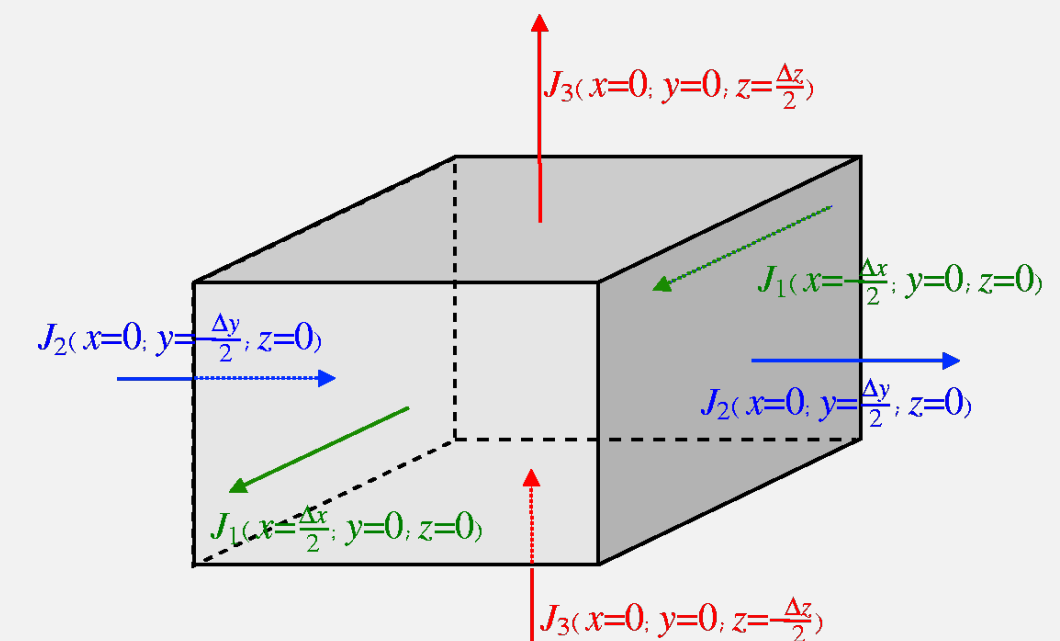
- $\frac{df(x)}{dx} = \lim_{\Delta x \rightarrow 0} \frac{f\left(x + \frac{\Delta x}{2}\right) - f\left(x - \frac{\Delta x}{2}\right)}{\Delta x} \approx \frac{f\left(x + \frac{\Delta x}{2}\right) - f\left(x - \frac{\Delta x}{2}\right)}{\Delta x}$

- Flux density

$$\bar{V}(x, y, z, t) = [J_1(x, y, z, t) \quad J_2(x, y, z, t) \quad J_3(x, y, z, t)]^T$$

- Net increase of material in the x-direction

$$\left\{ J_1\left(x + \frac{\Delta x}{2}, y, z, t\right) - J_1\left(x - \frac{\Delta x}{2}, y, z, t\right) \right\} \Delta y \Delta z \approx \frac{J_1\left(x + \frac{\Delta x}{2}, y, z, t\right) - J_1\left(x - \frac{\Delta x}{2}, y, z, t\right)}{\Delta x} \Delta x \Delta y \Delta z$$



Divergence (2/2)

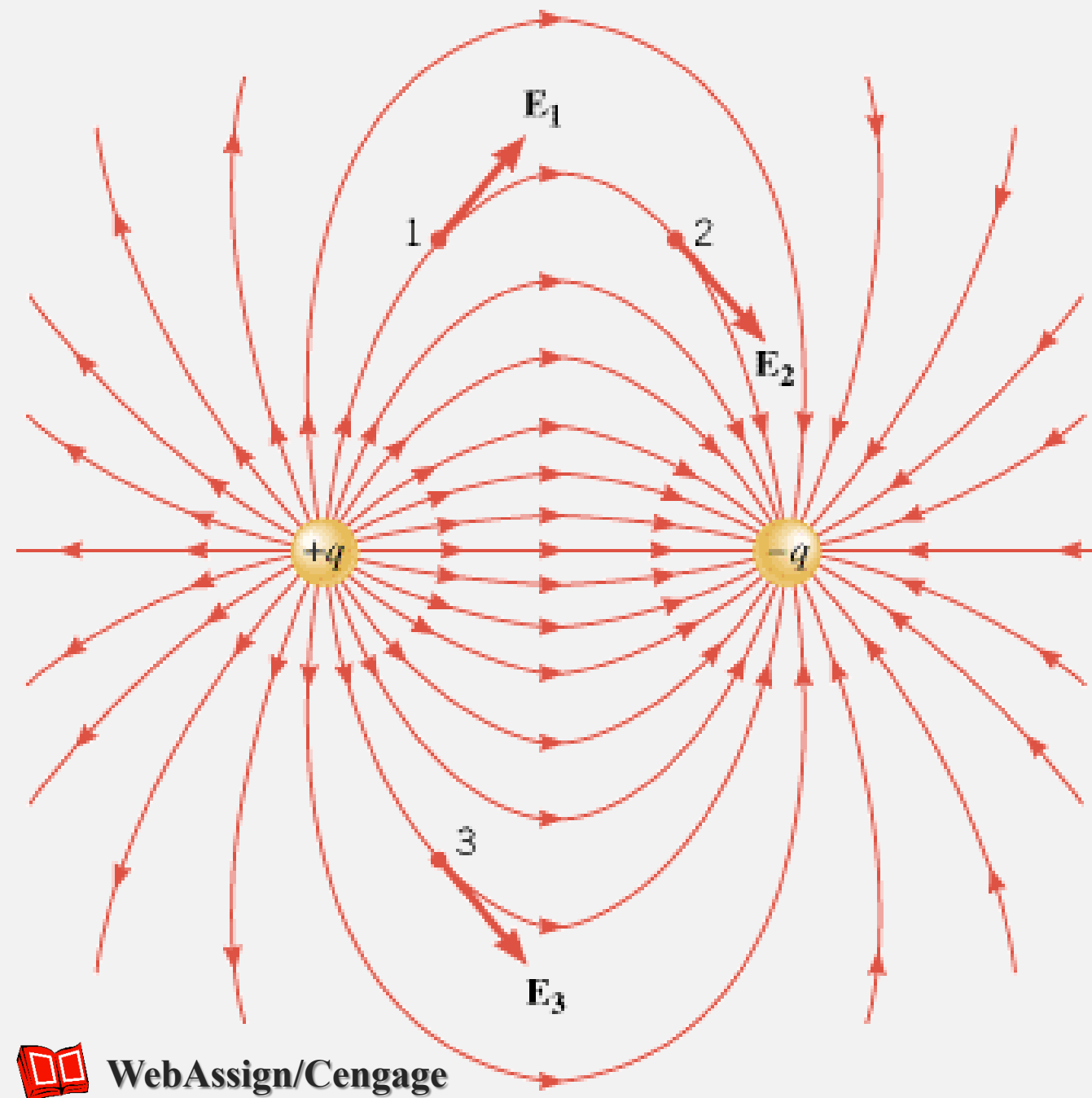
- Net increase of material per unit volume from the small cube
- $\text{div } \bar{V}(x, y, z, t) = \lim_{\Delta\tau \rightarrow 0} \left\{ \frac{\Delta J_1(x, y, z, t)}{\Delta x} + \frac{\Delta J_2(x, y, z, t)}{\Delta y} + \frac{\Delta J_3(x, y, z, t)}{\Delta z} \right\} =$
- $\frac{\partial J_1(x, y, z, t)}{\partial x} + \frac{\partial J_2(x, y, z, t)}{\partial y} + \frac{\partial J_3(x, y, z, t)}{\partial z}$
- With $\bar{V}(x, y, z, t) = \hat{x}V_x(x, y, z, t) + \hat{y}V_y(x, y, z, t) + \hat{z}V_z(x, y, z, t)$,

$$\text{div } \bar{V}(x, y, z, t) = \frac{\partial V_x(x, y, z, t)}{\partial x} + \frac{\partial V_y(x, y, z, t)}{\partial y} + \frac{\partial V_z(x, y, z, t)}{\partial z} =$$

$$\left(\hat{x} \frac{\partial}{\partial x} + \hat{y} \frac{\partial}{\partial y} + \hat{z} \frac{\partial}{\partial z} \right) \cdot \left(\hat{x}V_x(x, y, z, t) + \hat{y}V_y(x, y, z, t) + \hat{z}V_z(x, y, z, t) \right) =$$

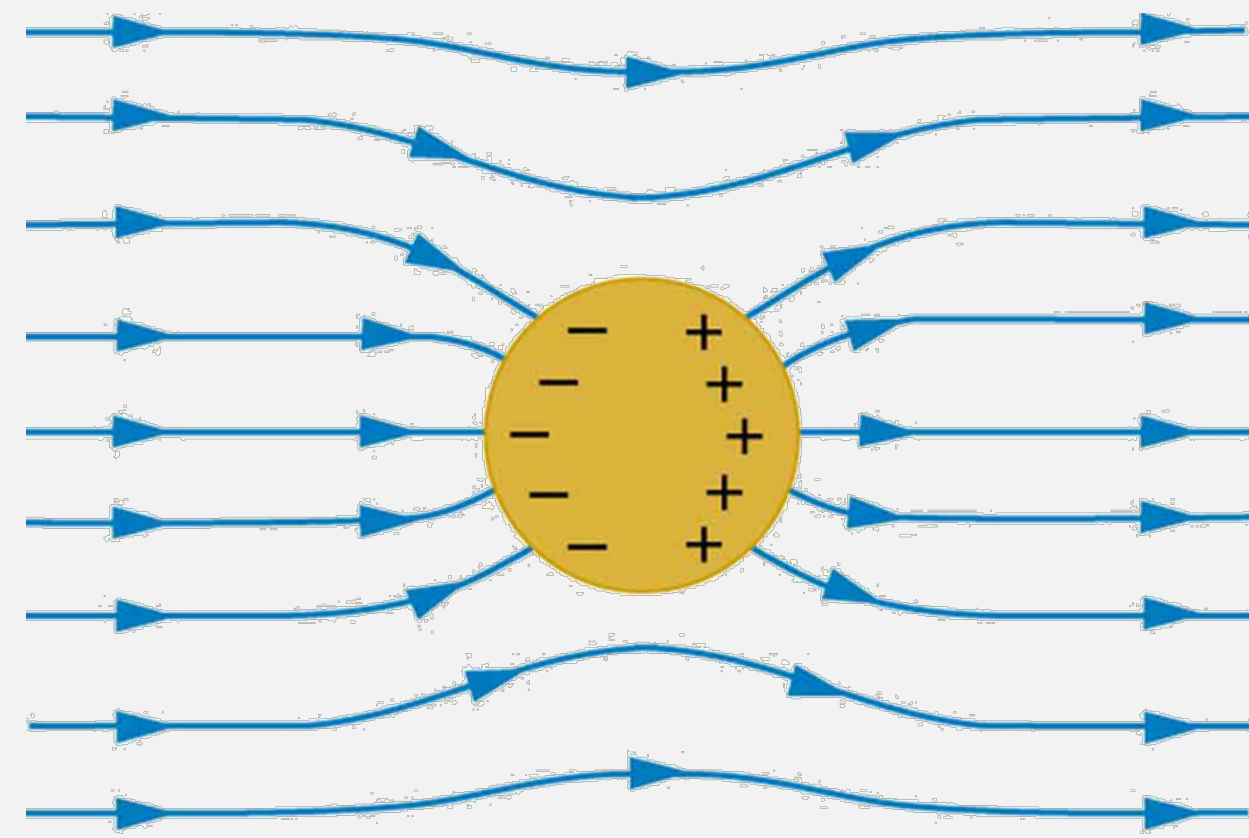
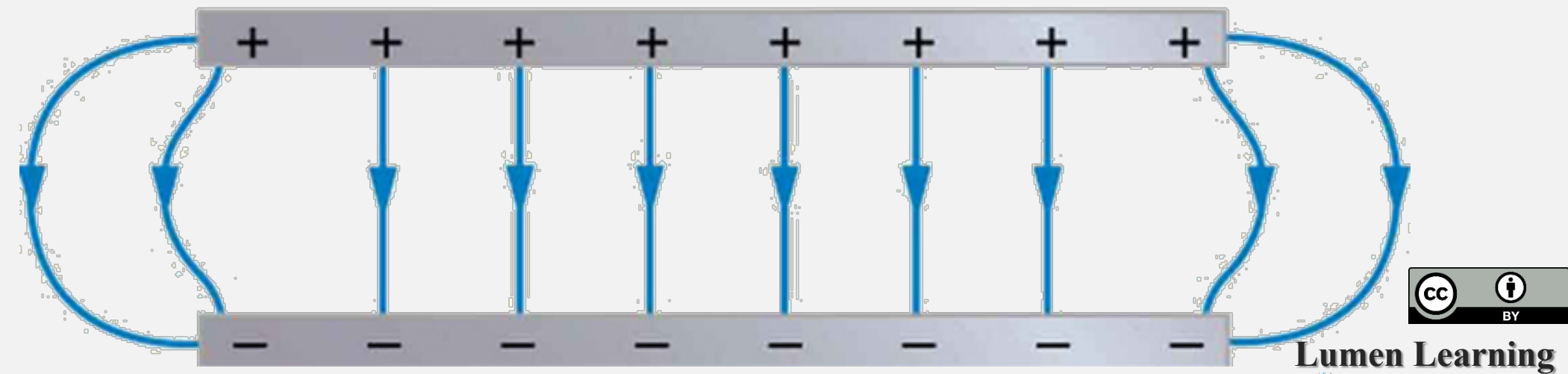
$$\nabla \cdot \bar{V}(x, y, z, t)$$

Sources and Sinks of Electric Field Lines



WebAssign/Cengage

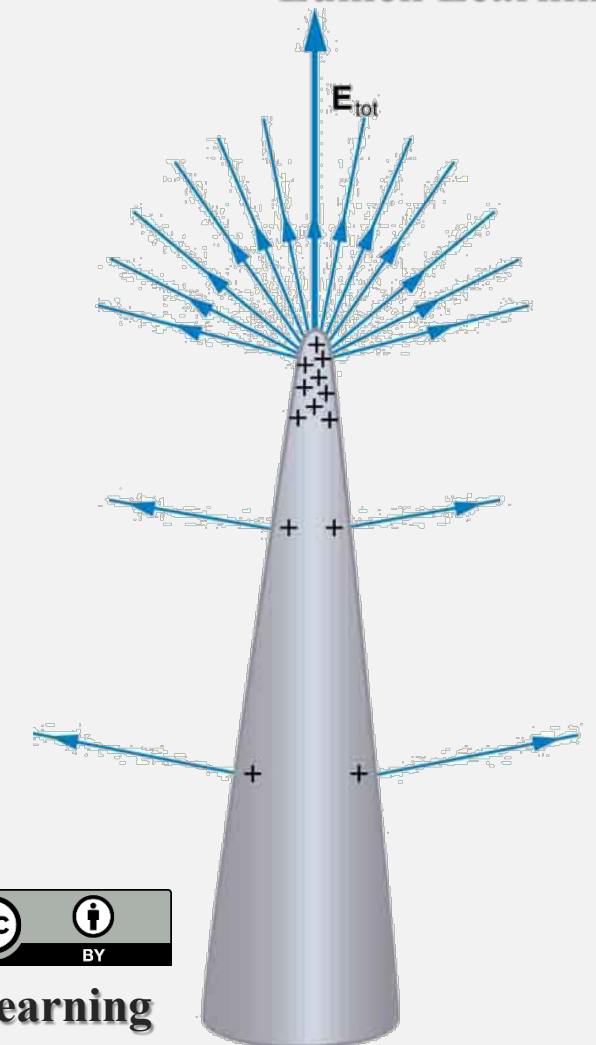
https://demo.webassign.net/ebooks/cj6demo/pc/c18/read/main/c18x18_7_ref.htm



Lumen Learning



Lumen Learning

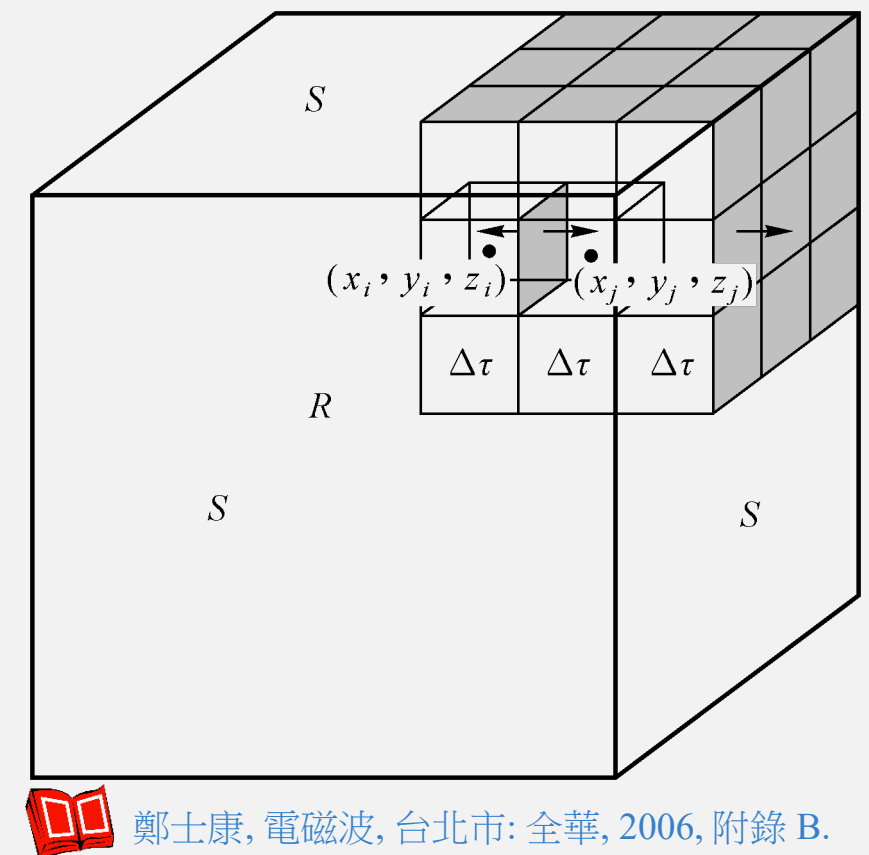


<https://courses.lumenlearning.com/physics/chapter/18-7-conductors-and-electric-fields-in-static-equilibrium/>

Divergence Theorem and Gauss's Law in Differential Form

- $\oint_S \bar{D}(x, y, z, t) \cdot \hat{n} da = \int_R \nabla \cdot \bar{D}(x, y, z, t) d\tau$
- $Q(t) = \int_R \rho_e(x, y, z, t) d\tau = \oint_S \bar{D}(x, y, z, t) \cdot \hat{n} da =$
 $\iint_{S_{yz}} D_x(x, y, z, t) n_x dy dz + \iint_{S_{zx}} D_y(x, y, z, t) n_y dz dx$
 $+ \iint_{S_{xy}} D_z(x, y, z, t) n_z dx dy = \int_R \nabla \cdot \bar{D}(x, y, z, t) d\tau$

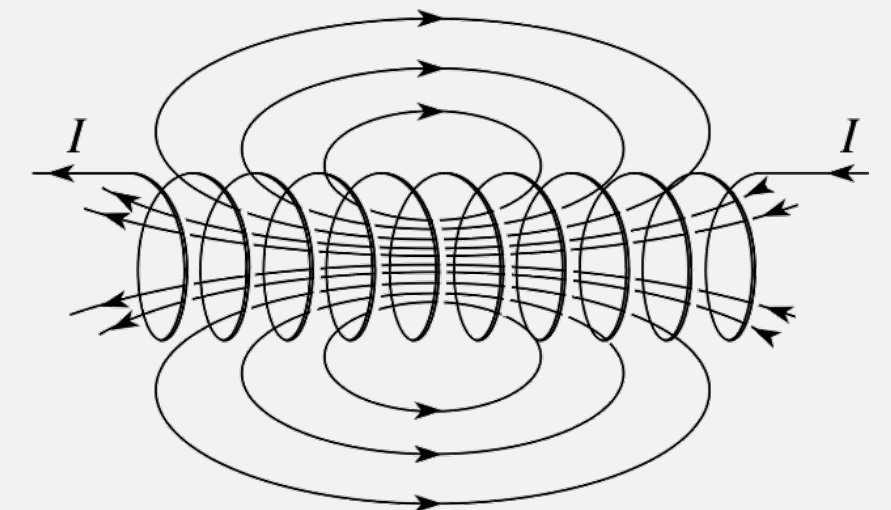
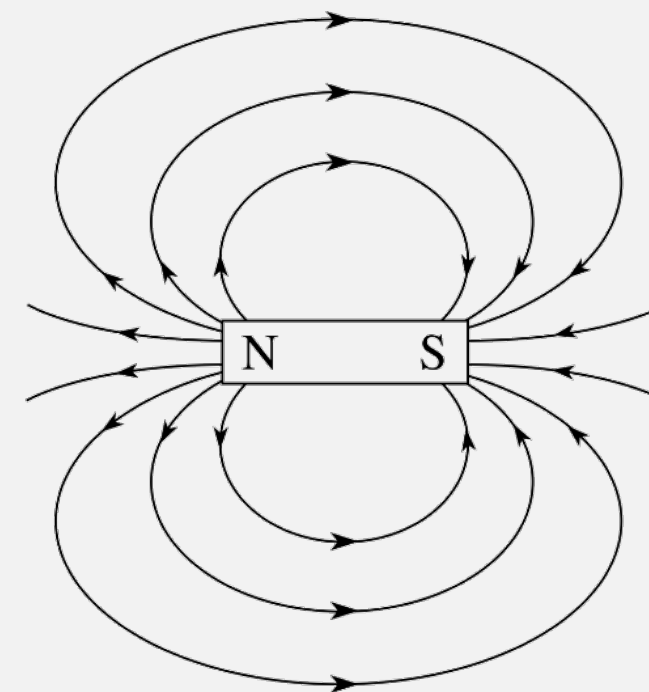
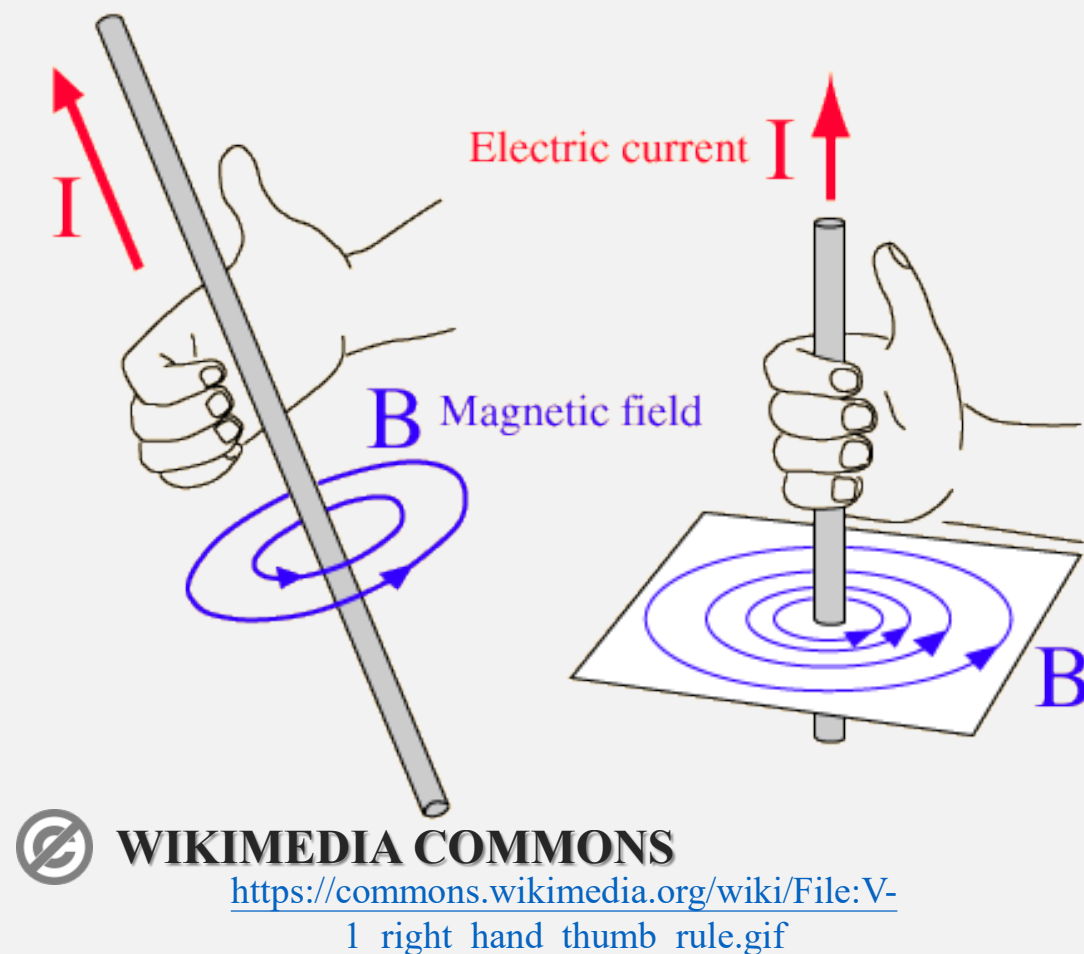
Let $R = \Delta\tau \rightarrow 0, \nabla \cdot \bar{D}(x, y, z, t) = \rho_e(x, y, z, t)$



鄭士康, 電磁波, 台北市: 全華, 2006, 附錄 B.

Gauss's Law for Magnetic Flux Density

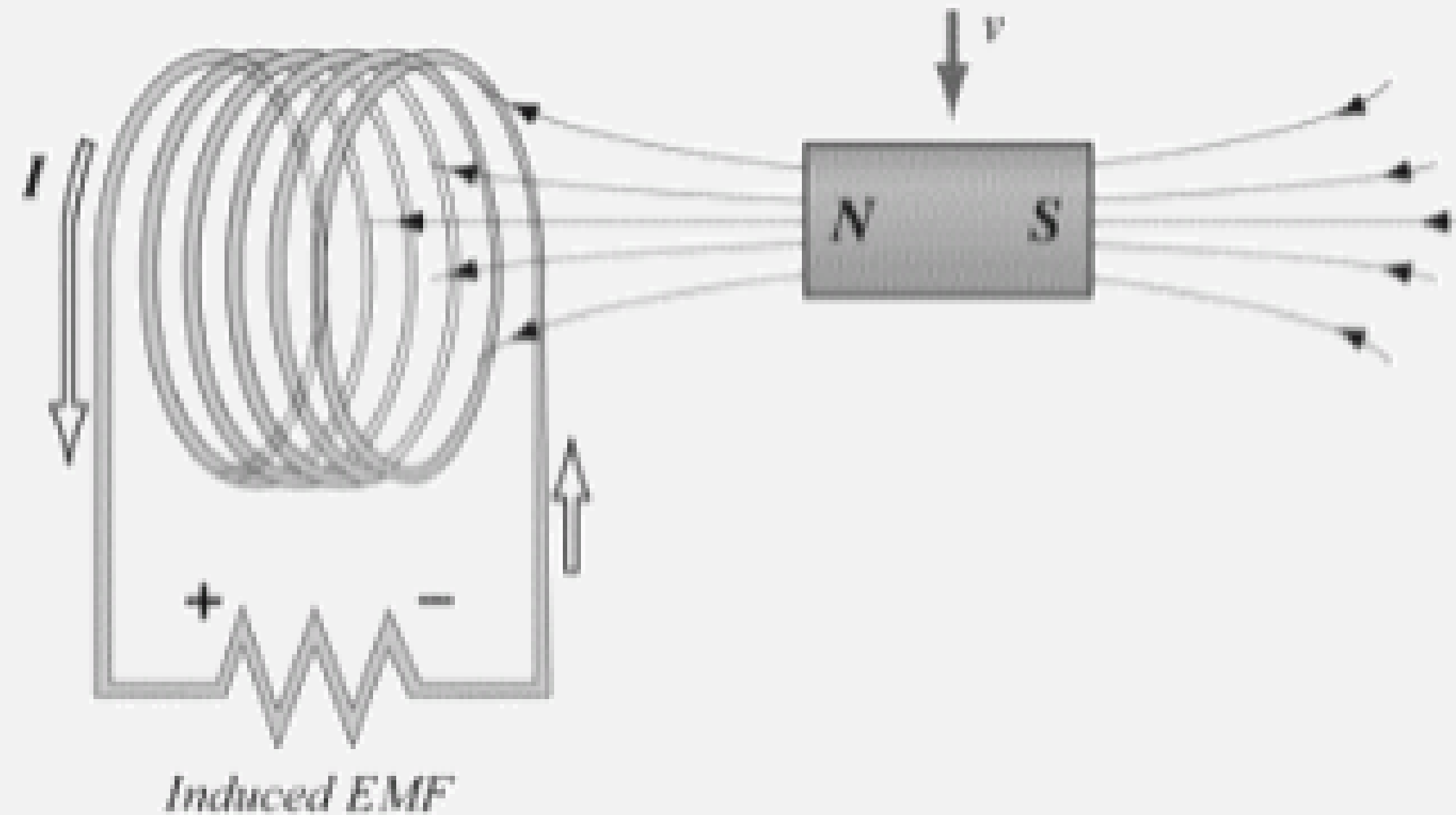
- $\oint_S \vec{B}(\vec{r}, t) \cdot \hat{n} da = 0$
- $\nabla \cdot \vec{B}(\vec{r}, t) = \frac{\partial B_x(\vec{r}, t)}{\partial x} + \frac{\partial B_y(\vec{r}, t)}{\partial y} + \frac{\partial B_z(\vec{r}, t)}{\partial z} = 0$



CBSE SCIENCE <https://cbscience.in/2017/09/27/magnetic-field-through-a-current-carrying-conductor/>

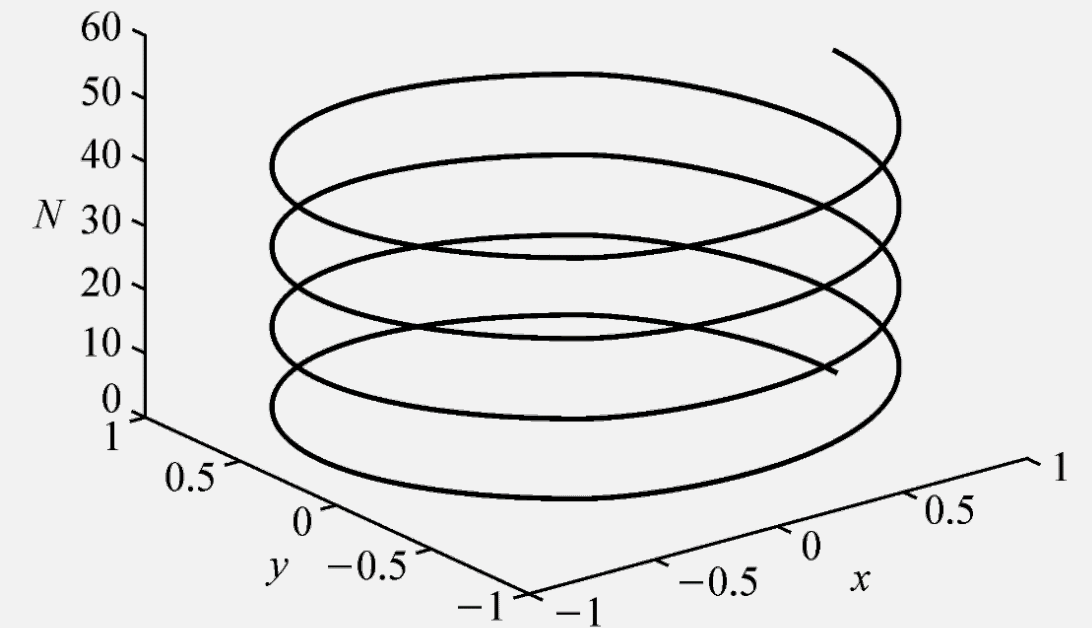
Faraday's Law

- $\varepsilon = -\frac{\Delta\Phi_M}{\Delta t}$



Line and Line Integral

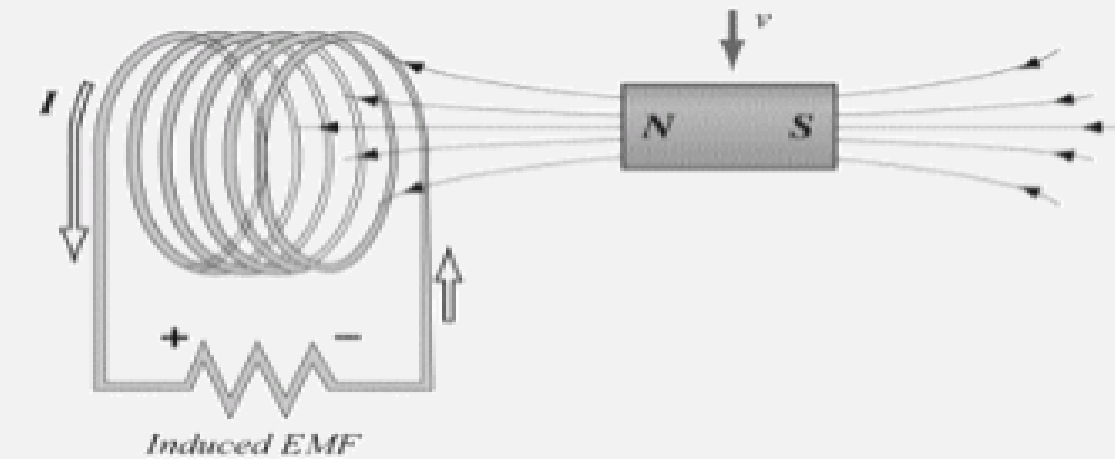
- Line $\vec{r} = \vec{r}(\sigma) = \hat{x}x(\sigma) + \hat{y}y(\sigma) + \hat{z}z(\sigma)$
- Example
- $\vec{r}(\sigma) = \hat{x} \cos \sigma + \hat{y} \sin \sigma + \hat{z} 2 \sigma, 0 \leq \sigma \leq 8\pi$
- $\hat{t} = (\hat{x}dx + \hat{y}dy + \hat{z}dz)/d\ell$
- $d\ell = \sqrt{(dx)^2 + (dy)^2 + (dz)^2} = \sqrt{\left(\frac{dx}{d\sigma}\right)^2 + \left(\frac{dy}{d\sigma}\right)^2 + \left(\frac{dz}{d\sigma}\right)^2} d\sigma$
- Line integral
- $\int_C \vec{F}(r, t) \cdot \hat{t} d\ell =$
 $\int_{C_\sigma} \vec{F}(r, t) \cdot \hat{t} \sqrt{\left(\frac{dx}{d\sigma}\right)^2 + \left(\frac{dy}{d\sigma}\right)^2 + \left(\frac{dz}{d\sigma}\right)^2} d\sigma =$
 $\int_{C_x} F_x(r, t) dx + \int_{C_y} F_y(r, t) dy + \int_{C_z} F_z(r, t) dz$



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Electric Field Strength and Faraday's Law

- $\bar{E}(\bar{r}, t) = \bar{F}(\bar{r}, t)/q$
- $\mathcal{E} = -\frac{\Delta\Phi_M}{\Delta t}$
- $\oint_C \bar{E}(\bar{r}, t) \cdot \hat{t} d\ell = -\frac{\partial}{\partial t} \int_S \bar{B}(\bar{r}, t) \cdot \hat{n} da$
- $\int_{C_x} E_x(\bar{r}, t) dx + \int_{C_y} E_y(\bar{r}, t) dy + \int_{C_z} E_z(\bar{r}, t) dz =$
 $-\frac{\partial}{\partial t} \iint_{S_{yz}} B_x(\bar{r}, t) n_x dy dz - \frac{\partial}{\partial t} \iint_{S_{zx}} B_y(\bar{r}, t) n_y dz dx$
 $-\frac{\partial}{\partial t} \iint_{S_{xy}} B_z(\bar{r}, t) n_z dx dy$

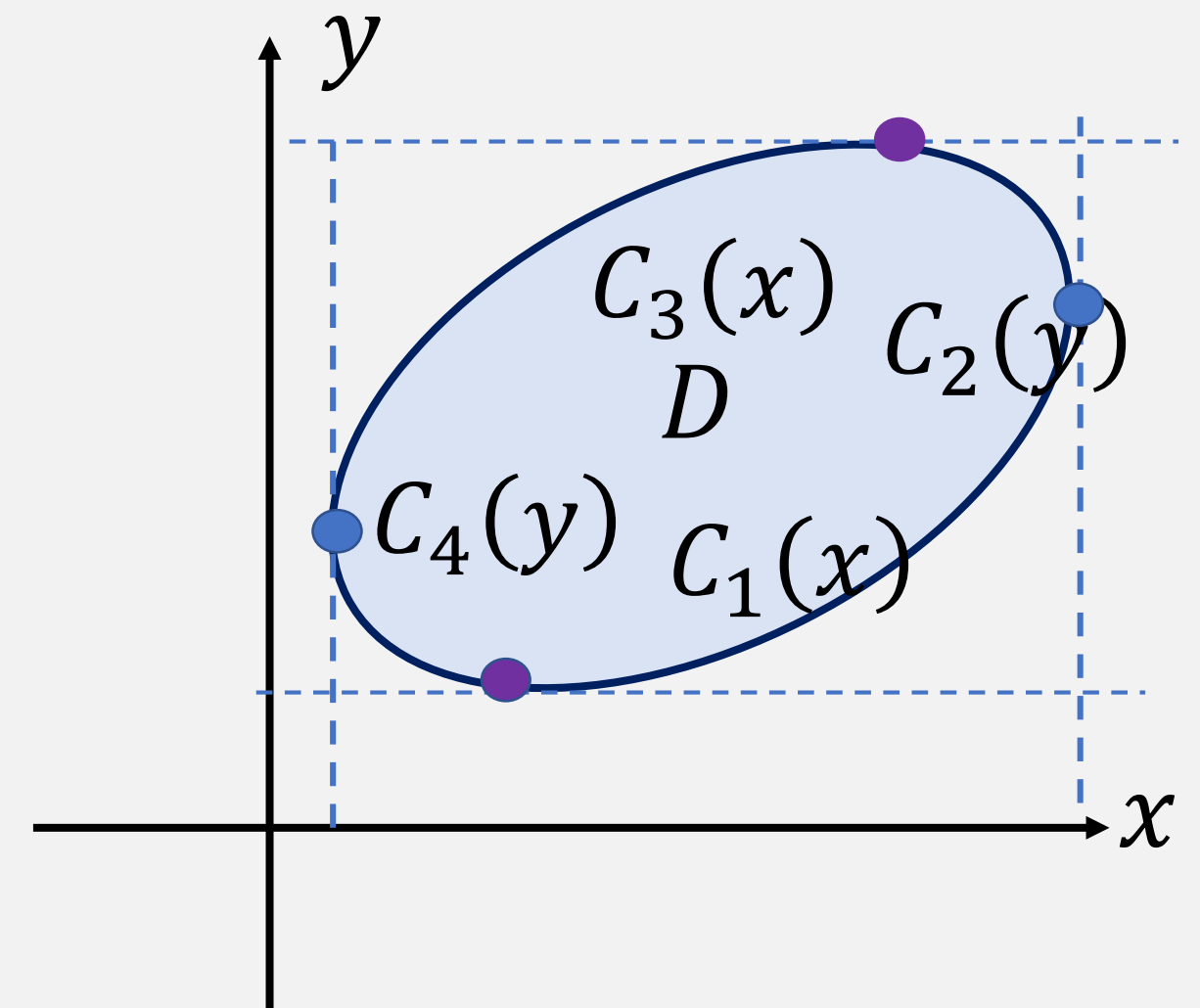


Suffolk University

<https://sites.suffolk.edu/kdshepard/2013/03/07/faradays-law/379a/>

Green's Theorem

- $\oint_C \bar{F}(x, y) \cdot \hat{t} d\ell = \int_C F_x(x, y) dx + \int_C F_y(x, y) dy$
- $\int_{C_x} F_x(x, y) dx = - \int_a^b \int_{C_1(x)}^{C_3(x)} \frac{\partial}{\partial y} F_x(x, y) dy dx$
- $\int_{C_y} F_y(x, y) dy = \int_c^d \int_{C_4(y)}^{C_2(y)} \frac{\partial}{\partial x} F_y(x, y) dx dy =$
- $\int_c^d \int_{C_2(y)}^{C_4(y)} \frac{\partial}{\partial x} F_y(x, y) dx dy$
- $\oint_C \bar{F}(x, y) \cdot \hat{t} d\ell =$
- $\int_{C_x} F_x(x, y) dx + \int_{C_y} F_y(x, y) dy =$
- $\int_D \left\{ \frac{\partial}{\partial x} F_y(x, y) - \frac{\partial}{\partial y} F_x(x, y) \right\} dx dy$



Stokes Theorem

- $\int_{C_x} F_x(\bar{r}, t) dx + \int_{C_y} F_y(\bar{r}, t) dy + \int_{C_z} F_z(\bar{r}, t) dz =$
- $\iint_{S_{yz}} \left\{ \frac{\partial}{\partial y} F_z(\bar{r}, t) - \frac{\partial}{\partial z} F_y(\bar{r}, t) \right\} n_x dydz + \iint_{S_{zx}} \left\{ \frac{\partial}{\partial z} F_x(\bar{r}, t) - \frac{\partial}{\partial x} F_z(\bar{r}, t) \right\} n_y dzdx + \iint_{S_{xy}} \left\{ \frac{\partial}{\partial x} F_y(\bar{r}, t) - \frac{\partial}{\partial y} F_x(\bar{r}, t) \right\} n_z dxdy$
- $\oint_C \bar{F}(\bar{r}, t) \cdot \hat{t} d\ell = \int_S \text{curl } \bar{F}(\bar{r}, t) \cdot \hat{n} da,$

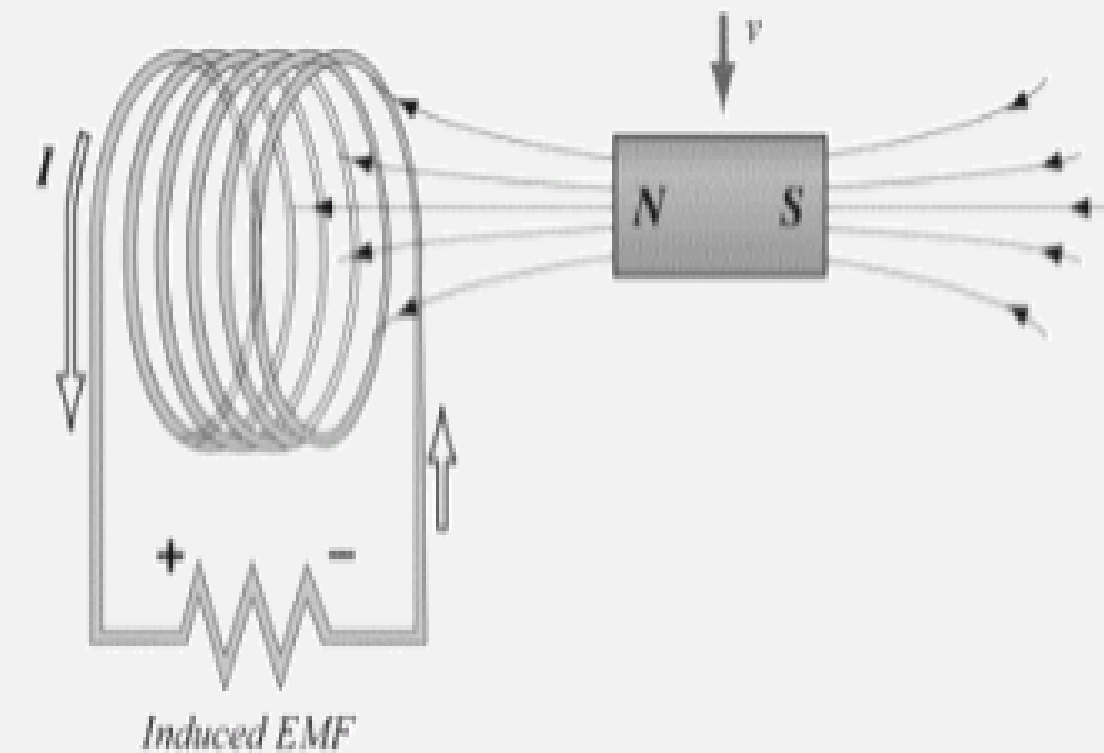
curl Operator

- $\oint_C \bar{F}(\bar{r}, t) \cdot \hat{t} d\ell = \int_S \text{curl } \bar{F}(\bar{r}, t) \cdot \hat{n} da,$
- $\text{curl } \bar{F}(\bar{r}, t) \cdot \hat{n} = \lim_{\Delta a \rightarrow 0} \frac{\oint_C \bar{F}(\bar{r}, t) \cdot \hat{t} d\ell}{\Delta a}$
- $\text{curl } \bar{F}(\bar{r}, t) = \hat{x} \left\{ \frac{\partial}{\partial y} F_z(\bar{r}, t) - \frac{\partial}{\partial z} F_y(\bar{r}, t) \right\} + \hat{y} \left\{ \frac{\partial}{\partial z} F_x(\bar{r}, t) - \frac{\partial}{\partial x} F_z(\bar{r}, t) \right\} + \hat{z} \left\{ \frac{\partial}{\partial x} F_y(\bar{r}, t) - \frac{\partial}{\partial y} F_x(\bar{r}, t) \right\} =$

$$\begin{vmatrix} \hat{x} & \hat{y} & \hat{z} \\ \frac{\partial}{\partial x} & \frac{\partial}{\partial y} & \frac{\partial}{\partial z} \\ F_x(\bar{r}, t) & F_y(\bar{r}, t) & F_z(\bar{r}, t) \end{vmatrix} = \nabla \times \bar{F}(\bar{r}, t)$$

Differential Form of Faraday's Law

- $\oint_C \vec{E}(\vec{r}, t) \cdot \hat{t} d\ell = -\frac{\partial}{\partial t} \int_S \vec{B}(\vec{r}, t) \cdot \hat{n} da$
- $\int_S \nabla \times \vec{E}(\vec{r}, t) \cdot \hat{n} da = -\frac{\partial}{\partial t} \int_S \vec{B}(\vec{r}, t) \cdot \hat{n} da$
- Let $s = \Delta s \rightarrow 0$
- $\hat{x} \cdot \nabla \times \vec{E}(\vec{r}, t) n_x \Delta s = -\frac{\partial}{\partial t} \hat{x} \cdot \frac{\partial}{\partial t} \vec{B}(\vec{r}, t) n_x \Delta s$
- $\hat{y} \cdot \nabla \times \vec{E}(\vec{r}, t) n_y \Delta s = -\frac{\partial}{\partial t} \hat{y} \cdot \frac{\partial}{\partial t} \vec{B}(\vec{r}, t) n_y \Delta s$
- $\hat{z} \cdot \nabla \times \vec{E}(\vec{r}, t) n_z \Delta s = -\frac{\partial}{\partial t} \hat{z} \cdot \frac{\partial}{\partial t} \vec{B}(\vec{r}, t) n_z \Delta s$
- $\nabla \times \vec{E}(\vec{r}, t) = -\frac{\partial}{\partial t} \vec{B}(\vec{r}, t)$

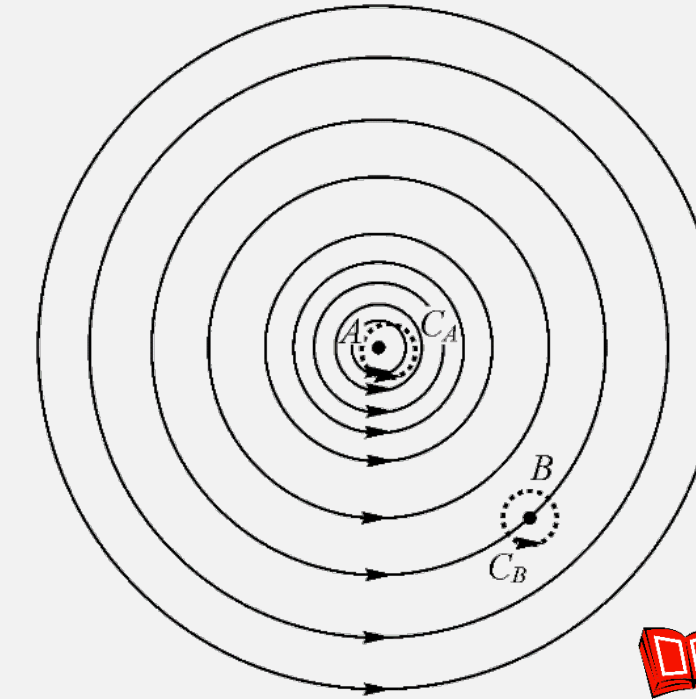


Suffolk University

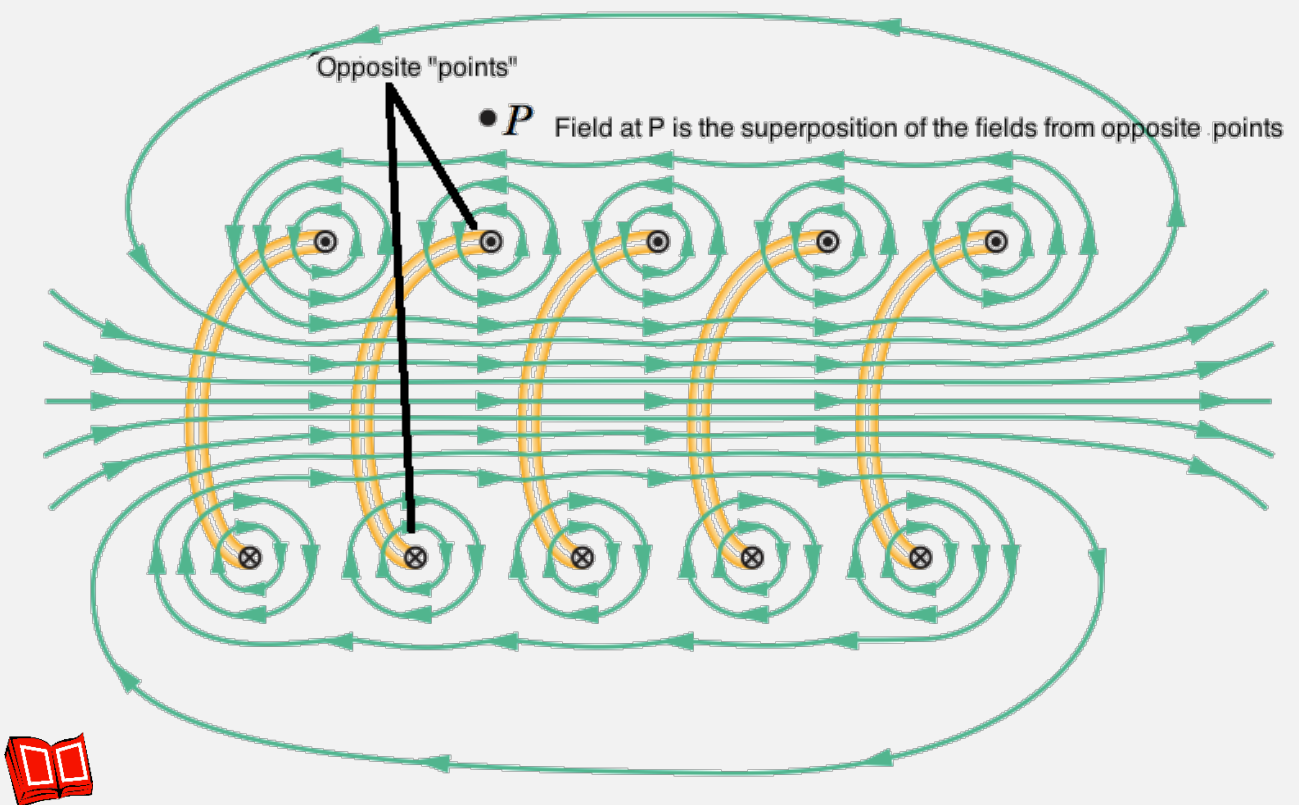
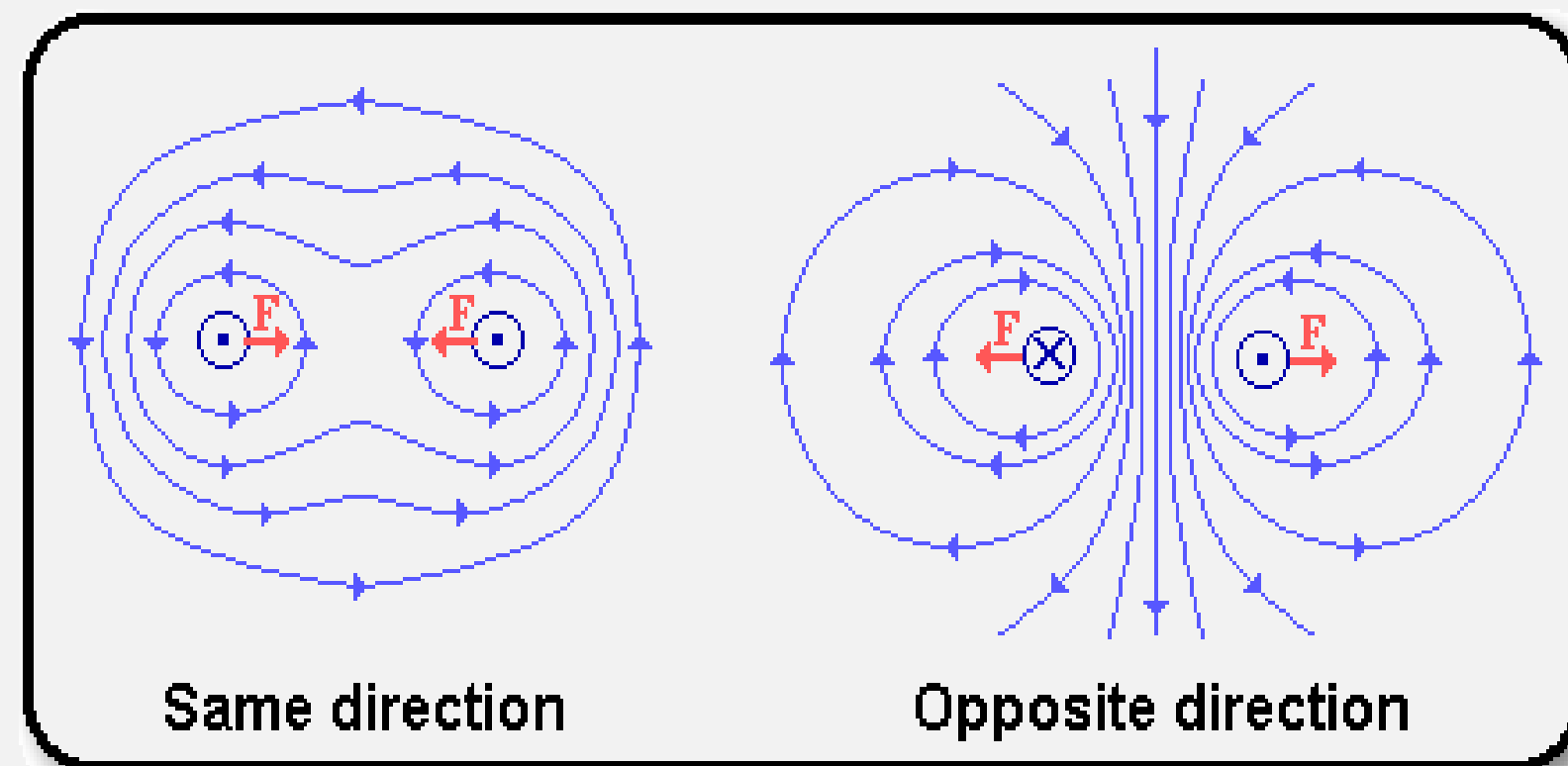
<https://sites.suffolk.edu/kdshepard/2013/03/07/faradays-law/379a/>

Ampère's Circuital Law

- $\oint_C \vec{H}(\vec{r}, t) \cdot \hat{t} d\ell = \int_S \vec{J}(\vec{r}, t) \cdot \hat{n} da$
- $\text{curl } \vec{H}(\vec{r}, t) = \nabla \times \vec{H}(\vec{r}, t) = \vec{J}(\vec{r}, t)$



 鄭士康, 電磁波, 台北市: 全華, 2006, 附錄 B.



Stock Exchange

<https://physics.stackexchange.com/questions/102312/what-are-the-limits-of-validity-for-the-magnetic-field-of-a-solenoid>

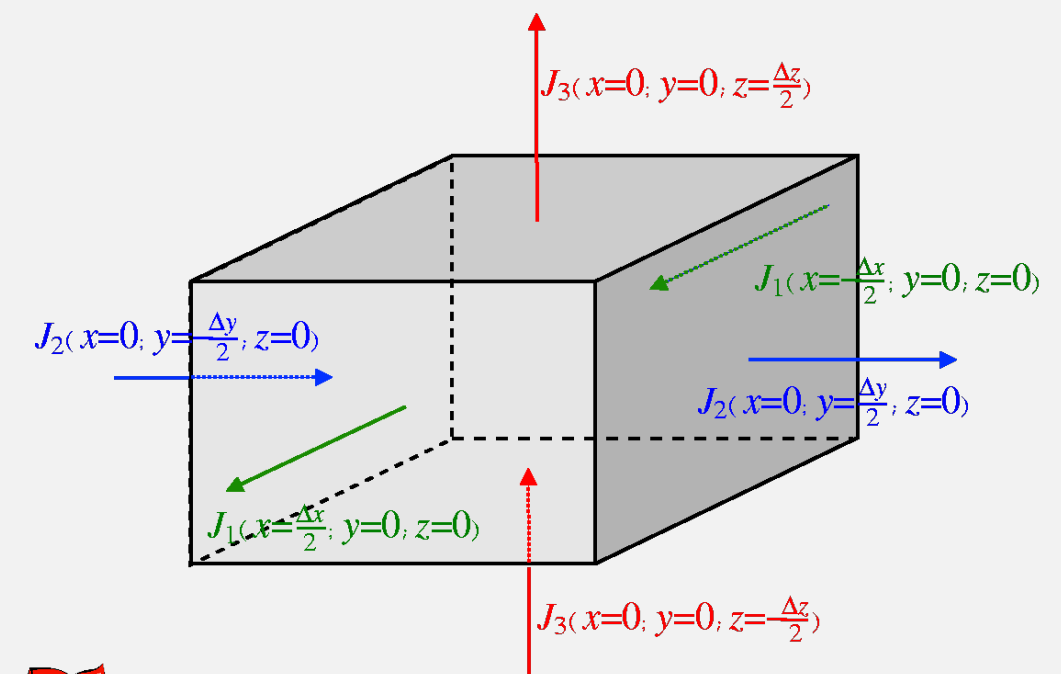


Magnetic field line due to two current carrying wires

<https://goprep.co/what-is-the-magnetic-field-midway-between-two-parallel-i-1nlo17>

Conservation of Charges

- $\oint_S \bar{J}(\bar{r}, t) \cdot \hat{n} da + \frac{\partial}{\partial t} Q(t) = 0$
- $\int_R \nabla \cdot \bar{J}(\bar{r}, t) d\tau + \frac{\partial}{\partial t} \int_R \rho_e(\bar{r}, t) d\tau = 0$
- $\nabla \cdot \bar{J}(\bar{r}, t) + \frac{\partial}{\partial t} \rho_e(\bar{r}, t) = 0$

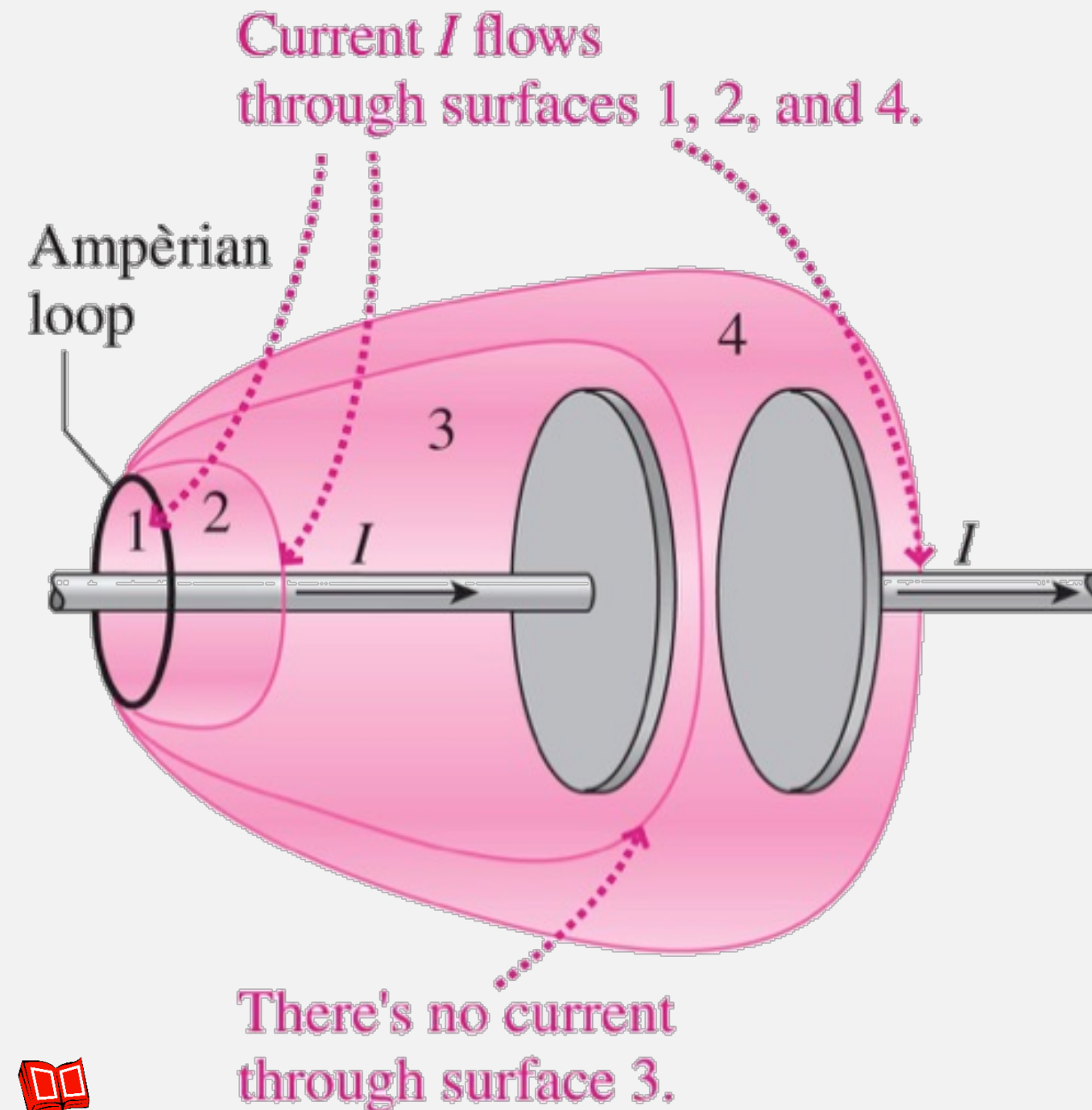


Massachusetts Institute of Technology

http://pruffle.mit.edu/3.016-2005/Lecture_16_web/node2.html

Inconsistency of Ampere's Circuital Law

- $\oint_C \vec{H}(\vec{r}, t) \cdot \hat{t} d\ell = \int_S \vec{J}(\vec{r}, t) \cdot \hat{n} da$
- $\nabla \times \vec{H}(\vec{r}, t) = \vec{J}(\vec{r}, t)$
- $\nabla \cdot \nabla \times \vec{H}(\vec{r}, t) = \nabla \cdot \vec{J}(\vec{r}, t)$
- $\nabla \cdot \nabla \times \vec{H}(\vec{r}, t) \equiv 0$
- $\nabla \cdot \vec{J}(\vec{r}, t) = 0 \neq -\frac{\partial}{\partial t} \rho_e(\vec{r}, t)$

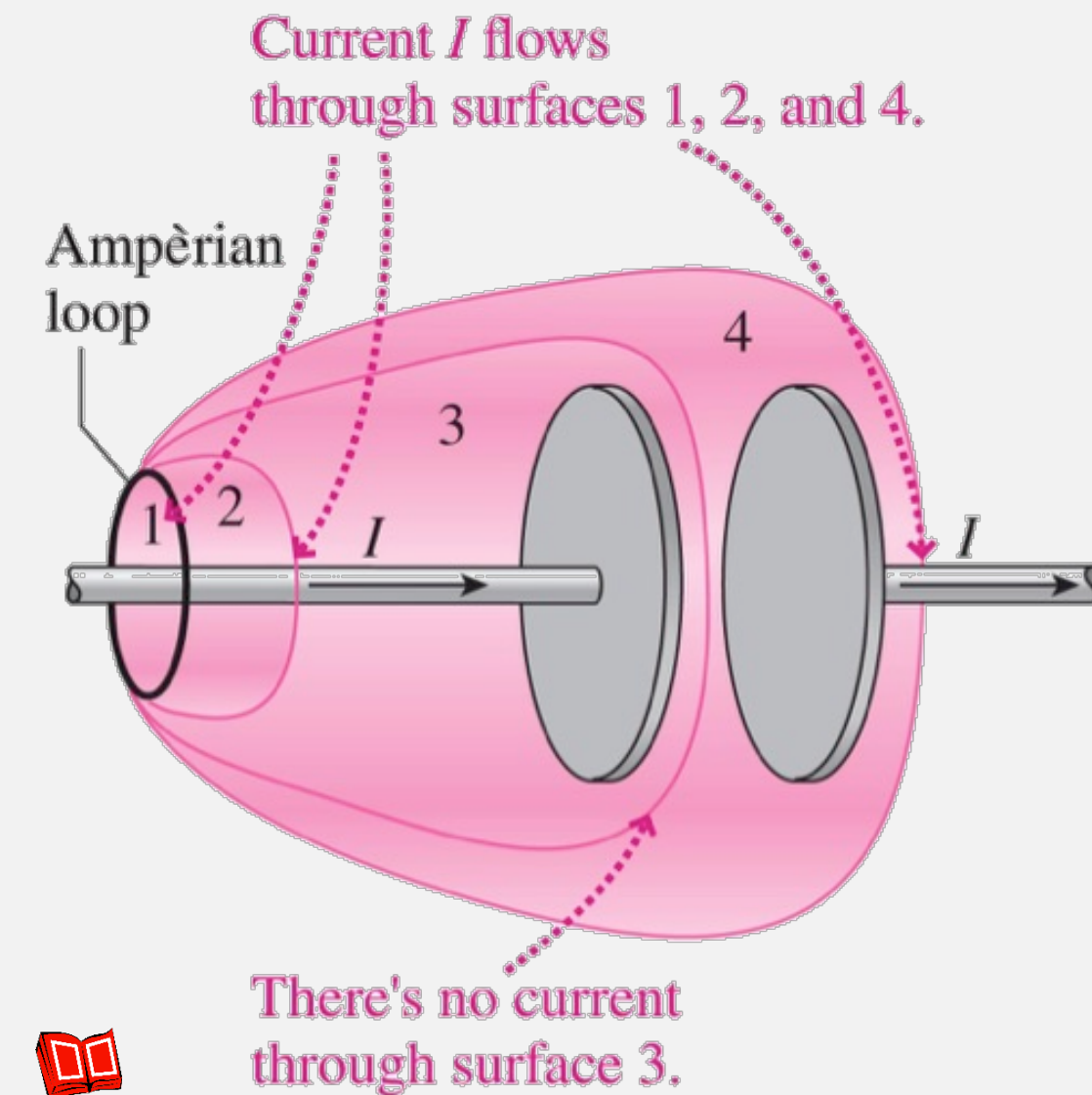


Physics Department, University of Louisville

https://www.physics.louisville.edu/cldavis/phys299/notes/mag_displacement.html

Maxwell's Displacement Current Density

- $\oint_C \vec{H}(\vec{r}, t) \cdot \hat{t} d\ell = \int_S \left\{ \vec{J}(\vec{r}, t) + \frac{\partial}{\partial t} \vec{D}(\vec{r}, t) \right\} \cdot \hat{n} da$
- $\nabla \times \vec{H}(\vec{r}, t) = \vec{J}(\vec{r}, t) + \frac{\partial}{\partial t} \vec{D}(\vec{r}, t)$
- $\nabla \cdot \nabla \times \vec{H}(\vec{r}, t) =$
 $0 = \nabla \cdot \vec{J}(\vec{r}, t) + \frac{\partial}{\partial t} \nabla \cdot \vec{D}(\vec{r}, t) =$
 $\nabla \cdot \vec{J}(\vec{r}, t) + \frac{\partial}{\partial t} \rho_e(\vec{r}, t)$



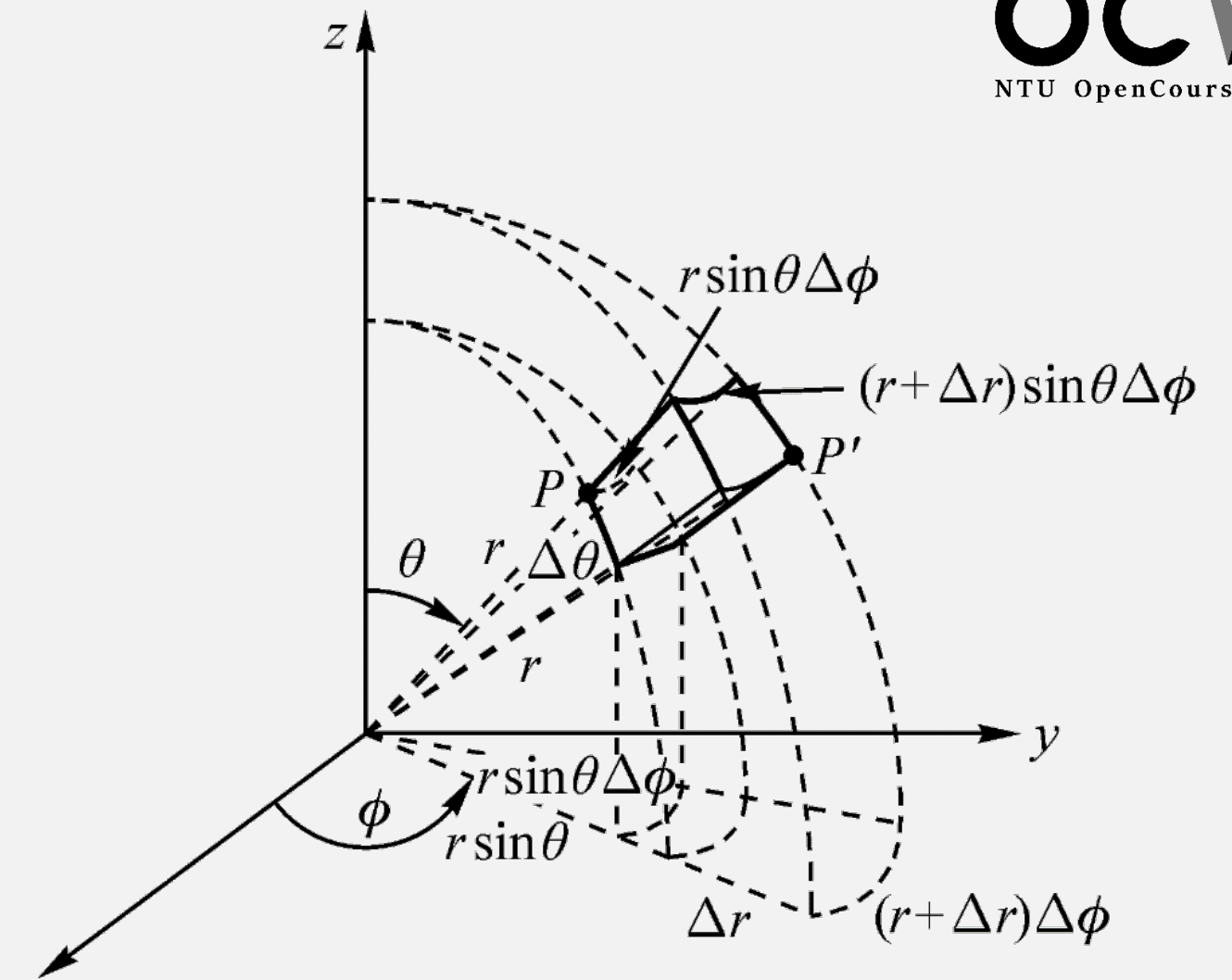
Physics Department, University of Louisville

https://www.physics.louisville.edu/cldavis/phys299/notes/mag_displacement.html

Permittivity of Free Space

- $\bar{D}(\bar{r}, t) = \epsilon_0 \bar{E}(\bar{r}, t)$
- $\oint_S \bar{D} \cdot \hat{n} da = q$
- Coulomb's law: $\bar{E} = K \frac{q}{r^2}$
- Spherical surface:

$$\bar{r}(\theta, \phi) = \hat{x} r \sin \theta \cos \phi + \hat{y} r \sin \theta \sin \phi + \hat{z} r \cos \theta$$
- Differential solid angle: $d\Omega = \sin \theta d\theta d\phi$
- $q = \int_0^{2\pi} \int_0^\pi \epsilon_0 K \frac{q}{r^2} r^2 \sin \theta d\theta d\phi = 4\pi \epsilon_0 K q$
- $K = \frac{1}{4\pi \epsilon_0}$

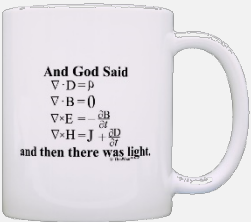


鄭士康, 電磁波, 台北市: 全華, 2006, 附錄 B.

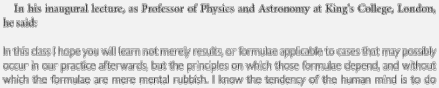
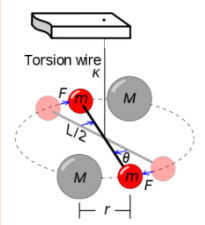
Permeability of Free Space

- $\bar{B}(\bar{r}, t) = \mu_0 \bar{H}(\bar{r}, t)$

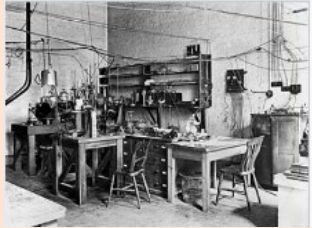
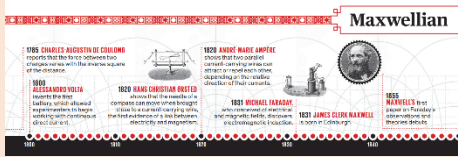
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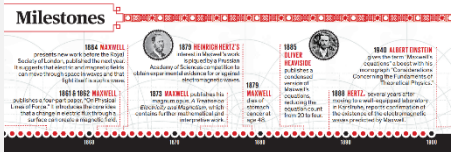
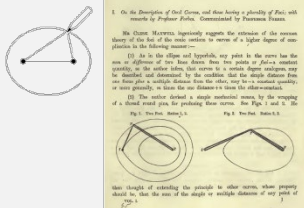
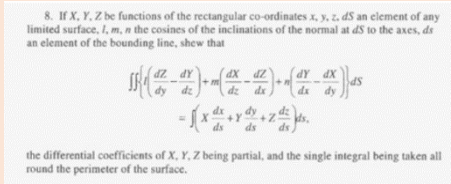
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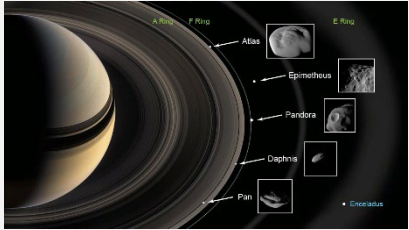
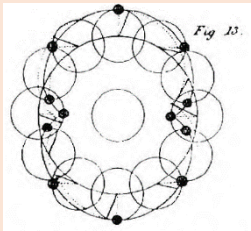
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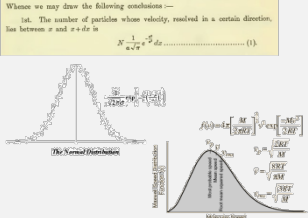
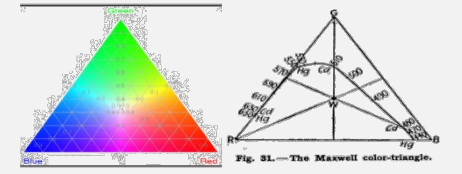
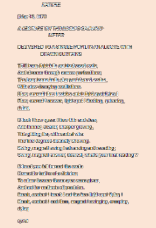
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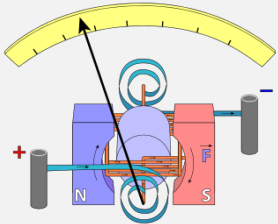
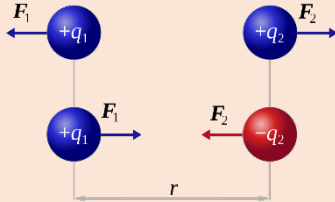
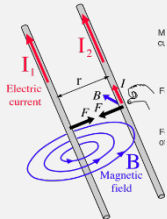
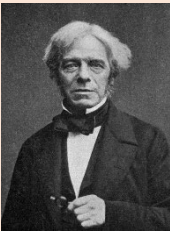
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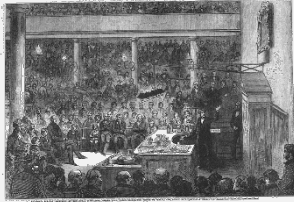
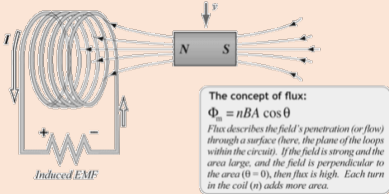
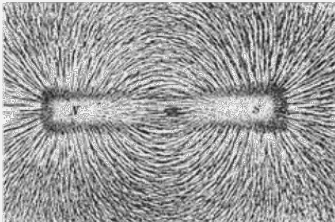
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18、46、48、52			Suffolk University / Discovering The World / Faraday's Law , https://sites.suffolk.edu/kdshepard/2013/03/07/faradays-law/379a/ , 本網站係以著作權法第 46、52、65 條合理使用本件作品，2022/01/27 瀏覽。
18-19、22、35			Google Books / Practical Physics / Newton Henry Black , https://books.google.com.tw/books?id=BT8AAAAAYAAJ&pg=PA242&redir_esc=y#v=onepage&q&f=false , 本網站係以著作權法第 46、52、65 條合理使用本件作品，2022/01/27 瀏覽。
19	... by line of magnetic force, be either in curved or straight lines.		Michael Faraday, Experimental Researches in Electricity, vol. III, London: Bernard Quaritch, 1855, ¶ 2149. , 本網站係以著作權法第 46、52、65 條合理使用本件作品，2022/01/27 瀏覽。

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21	By a physical analogy of quantities by operations on numbers.		James Clerk Maxwell, On Faraday’s Lines of Force, Transactions of the Cambridge Philosophical Society, vol. x, part 1, 1855. p. 155.，本網站係以著作權法第 46、52、65 條合理使用本件作品。
22	“[T]his curve will indicate the direction of that on that account <i>a line of force</i> [Maxwell’s emphasis].”		James Clerk Maxwell, On Faraday’s Lines of Force, Transactions of the Cambridge Philosophical Society, vol. x, part 1, 1855. p. 158.，本網站係以著作權法第 46、52、65 條合理使用本件作品。
22	“[C]onsider these curves not merely lines, by the motion of the fluid in these tubes.”		James Clerk Maxwell, On Faraday’s Lines of Force, Transactions of the Cambridge Philosophical Society, vol. x, part 1, 1855. p. 159.，本網站係以著作權法第 46、52、65 條合理使用本件作品。
23	Three equations of Magnetic Force (B) “ Electric Currents (C) “ Electromotive Force (D) “ Electric Elasticity (E) “ Electric Resistance (F) “ Total Currents (A) One equation of Free Electricity (G) “ Continuity (H)		James Clerk Maxwell, “A dynamic theory of the electromagnetic field,” Royal Society Transactions, vol. 155, 1865, pp. 459-512.，本網站係以著作權法第 46、52、65 條合理使用本件作品。

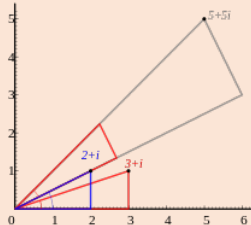
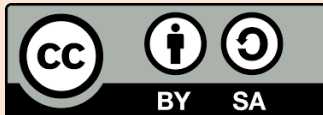
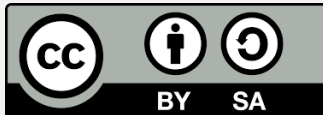
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23	For Electromagnetic Momentum F G H " Magnetic Intensity a b c " Electromotive Force P Q R " Current due to true conduction p q r " Electric Displacement f g h " Total Current (including variation of displacement) p' q' r' " Quantity of free Electricity e " Electric Potential Ψ		James Clerk Maxwell, “A dynamic theory of the electromagnetic field,” Royal Society Transactions, vol. 155, 1865, pp. 459-512.，本網站係以著作權法第 46、52、65 條合理使用本件作品。
24	$\left. \begin{aligned} p' &= p + \frac{df}{dt}, \\ q' &= q + \frac{dg}{dt}, \\ r' &= r + \frac{dh}{dt}, \end{aligned} \right\}$		James Clerk Maxwell, “A dynamic theory of the electromagnetic field,” Royal Society Transactions, vol. 155, 1865, pp. 459-512.，本網站係以著作權法第 46、52、65 條合理使用本件作品。
24	$\left. \begin{aligned} \mu\alpha &= \frac{dH}{dy} - \frac{dG}{dz}, \\ \mu\beta &= \frac{dF}{dz} - \frac{dH}{dx}, \\ \mu\gamma &= \frac{dG}{dx} - \frac{dF}{dy}. \end{aligned} \right\}$		James Clerk Maxwell, “A dynamic theory of the electromagnetic field,” Royal Society Transactions, vol. 155, 1865, pp. 459-512.，本網站係以著作權法第 46、52、65 條合理使用本件作品。
24	$\left. \begin{aligned} \frac{d\gamma}{dy} - \frac{d\beta}{dz} &= 4\pi p', \\ \frac{d\alpha}{dz} - \frac{d\gamma}{dx} &= 4\pi q', \\ \frac{d\beta}{dx} - \frac{d\alpha}{dy} &= 4\pi r'. \end{aligned} \right\}$		James Clerk Maxwell, “A dynamic theory of the electromagnetic field,” Royal Society Transactions, vol. 155, 1865, pp. 459-512.，本網站係以著作權法第 46、52、65 條合理使用本件作品。

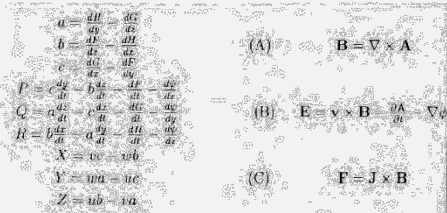
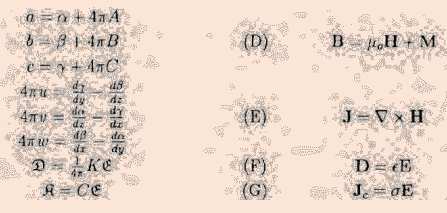
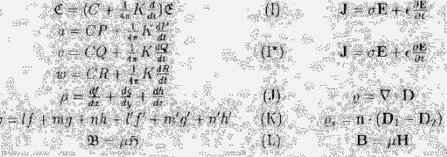
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25	$\left. \begin{aligned} P &= \mu \left(\gamma \frac{dy}{dt} - \beta \frac{dz}{dt} \right) - \frac{dF}{dt} - \frac{d\Psi}{dx}, \\ Q &= \mu \left(\alpha \frac{dz}{dt} - \gamma \frac{dx}{dt} \right) - \frac{dG}{dt} - \frac{d\Psi}{dy}, \\ R &= \mu \left(\beta \frac{dx}{dt} - \alpha \frac{dy}{dt} \right) - \frac{dH}{dt} - \frac{d\Psi}{dz}. \end{aligned} \right\}$		James Clerk Maxwell, “A dynamic theory of the electromagnetic field,” Royal Society Transactions, vol. 155, 1865, pp. 459-512.，本網站係以著作權法第 46、52、65 條合理使用本件作品。
25	$\left. \begin{aligned} P &= kf, \\ Q &= kg, \\ R &= kh. \end{aligned} \right\}$		James Clerk Maxwell, “A dynamic theory of the electromagnetic field,” Royal Society Transactions, vol. 155, 1865, pp. 459-512.，本網站係以著作權法第 46、52、65 條合理使用本件作品。
25	$\left. \begin{aligned} P &= -\varepsilon p, \\ Q &= -\varepsilon q, \\ R &= -\varepsilon r. \end{aligned} \right\}$		James Clerk Maxwell, “A dynamic theory of the electromagnetic field,” Royal Society Transactions, vol. 155, 1865, pp. 459-512.，本網站係以著作權法第 46、52、65 條合理使用本件作品。
25	$e + \frac{df}{dx} + \frac{dg}{dy} + \frac{dh}{dz} = 0.$		James Clerk Maxwell, “A dynamic theory of the electromagnetic field,” Royal Society Transactions, vol. 155, 1865, pp. 459-512.，本網站係以著作權法第 46、52、65 條合理使用本件作品。

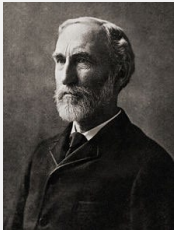
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25	$\frac{de}{dt}+\frac{dp}{dx}+\frac{dq}{dy}+\frac{dr}{dz}=0$		James Clerk Maxwell, “A dynamic theory of the electromagnetic field,” Royal Society Transactions, vol. 155, 1865, pp. 459-512.，本網站係以著作權法第 46、52、65 條合理使用本件作品。																									
26			WIKIMEDIA COMMONS / File:Complex multi.svg / 2 August 2015 / IkamusumeFan ， https://commons.wikimedia.org/wiki/File:Complex_multi.svg ， CC BY-SA 4.0 ，2022/01/27 瀏覽。																									
26	<table><tr><td>×</td><td>1</td><td>i</td><td>j</td><td>k</td></tr><tr><td>1</td><td>1</td><td>i</td><td>j</td><td>k</td></tr><tr><td>i</td><td>i</td><td>-1</td><td>k</td><td>-j</td></tr><tr><td>j</td><td>j</td><td>-k</td><td>-1</td><td>i</td></tr><tr><td>k</td><td>k</td><td>j</td><td>-i</td><td>-1</td></tr></table>	×	1	i	j	k	1	1	i	j	k	i	i	-1	k	-j	j	j	-k	-1	i	k	k	j	-i	-1		WIKIPEDIA / Quaternion， https://en.wikipedia.org/wiki/Quaternion ， CC BY-SA 3.0 ，2022/01/27 瀏覽。
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26			WIKIMEDIA COMMONS / File:William Rowan Hamilton portrait oval combined.png / 30 August 2010, 11:19 (UTC)， https://commons.wikimedia.org/wiki/File:William_Rowan_Hamilton_portrait_oval_combined.png ， 本網站係以著作權法第 46、52、65 條合理使用本件作品，2022/01/27 瀏覽。																									

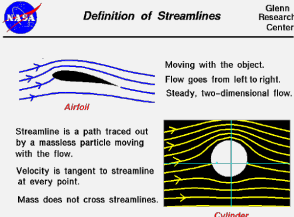
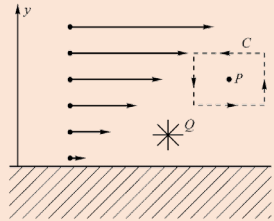
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28			Tapan K. Sarkar, etc., History of Wireless, New Jersey: John Wiley & Sons, Inc., 2006, Fig. 3.11, p. 185. / James Clerk Maxwell, A Treaties on Electricity and Magnetism, Oxford: Clarendon Press, 1873, pp. 480-493. , 本網站係以著作權法第 46、52、65 條合理使用本件作品。
29			Tapan K. Sarkar, etc., History of Wireless, New Jersey: John Wiley & Sons, Inc., 2006, Fig. 3.11, p. 185. / James Clerk Maxwell, A Treaties on Electricity and Magnetism, Oxford: Clarendon Press, 1873, pp. 480-493. , 本網站係以著作權法第 46、52、65 條合理使用本件作品。
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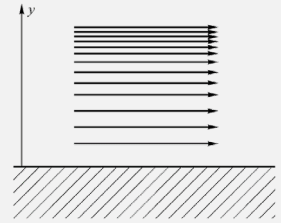
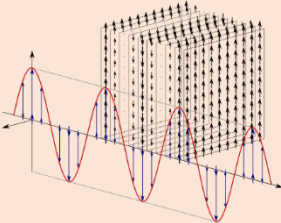
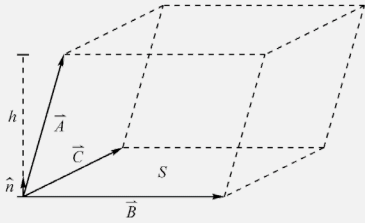
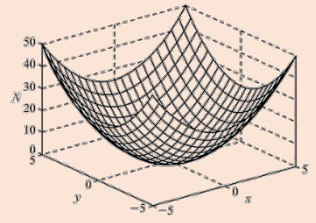
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30	$AB = AB \cos \theta$		Oliver Heaviside, <i>Electromagnetic Theory</i> , New York: Dover Publications, Inc., 1950, p. 38, ¶ 107. , 本網站係以著作權法第 46、52、65 條合理使用本件作品。
30	$C = VAB, C = AB \sin \theta$		Oliver Heaviside, <i>Electromagnetic Theory</i> , New York: Dover Publications, Inc., 1950, p. 40, ¶ 111. , 本網站係以著作權法第 46、52、65 條合理使用本件作品。
31	“used,..., to indicate the often-occurring operation of a line-integral in a closed circuit”		Oliver Heaviside, <i>Electromagnetic Theory</i> , New York: Dover Publications, Inc., 1950, p. 9, ¶ 33. , 本網站係以著作權法第 46、52、65 條合理使用本件作品。

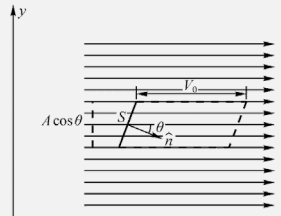
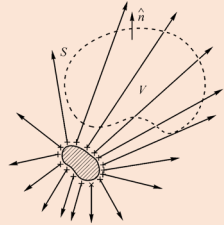
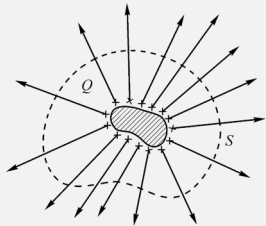
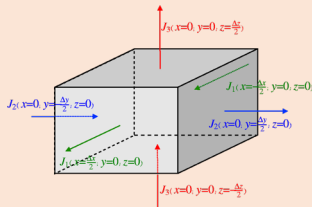
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31	“The divergence of any flux is the amount of the flux leaving the unit volume.”		Oliver Heaviside, Electromagnetic Theory, New York: Dover Publications, Inc., 1950, p. 13, ¶ 51. , 本網站係以著作權法第 46、52、65 條合理使用本件作品。
35			NASA , https://www.grc.nasa.gov/WWW/k-12/VirtualAero/BottleRocket/airplane/stream.html , 本網站係以著作權法第 46、52、65 條合理使用本件作品，2022/01/27 瀏覽。
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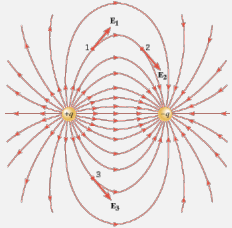
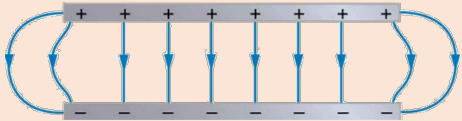
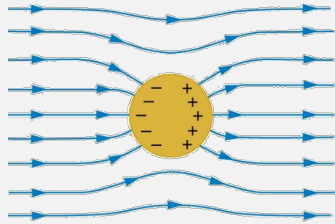
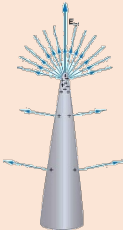
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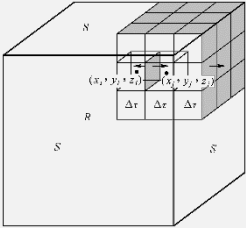
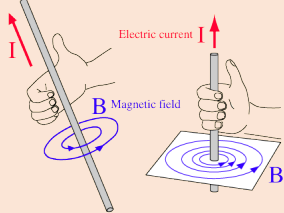
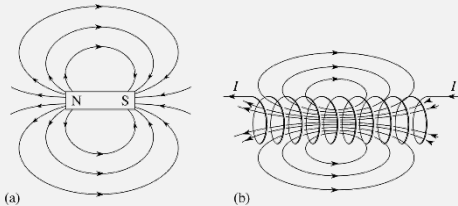
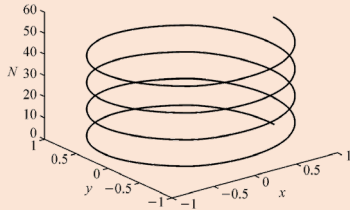
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41、54			Massachusetts Institute of Technology / The Divergence Theorem , http://pruffle.mit.edu/3.016-2005/Lecture_16_web/node2.html , 本網站係以著作權法第 46、52、65 條合理使用本件作品，2022/01/27 瀏覽。

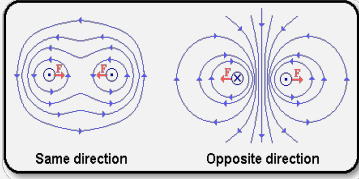
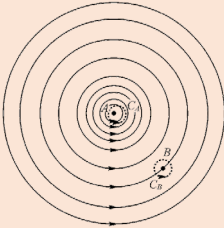
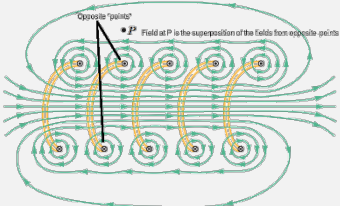
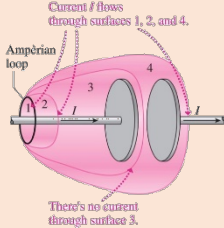
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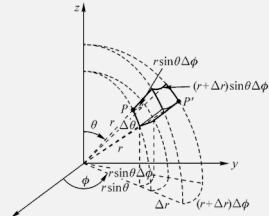
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45			CBSE SCIENCE / Learning Resource for CBSE students , https://cbsescience.in/2017/09/27/magnetic-field-through-a-current-carrying-conductor/ , 本網站係以著作權法第 46、52、65 條合理使用本件作品，2022/01/27 瀏覽。
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53	 Fig. Magnetic field line due to two current carrying wires		goprep / What is the magne , https://goprep.co/what-is-the-magnetic-field-midway-between-two-parallel-i-1nlo17 , 本網站係以著作權法第 46、52、65 條合理使用本件作品，2022/01/27 瀏覽。
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